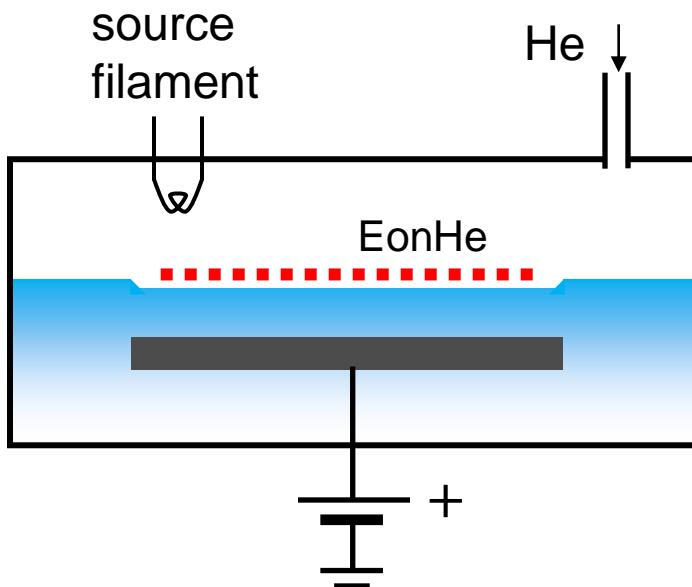
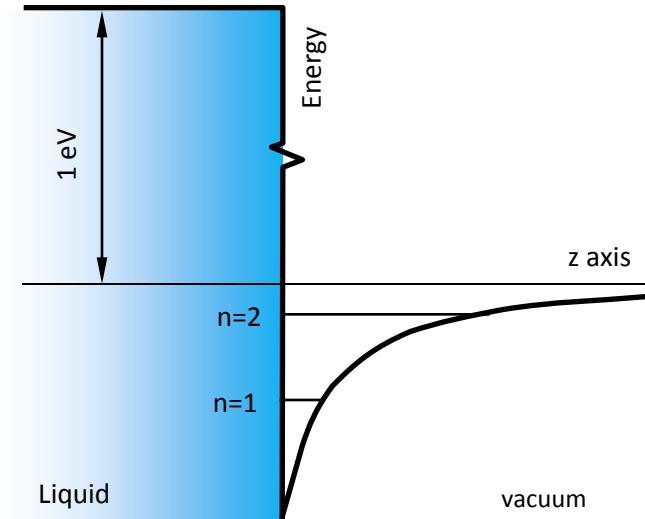
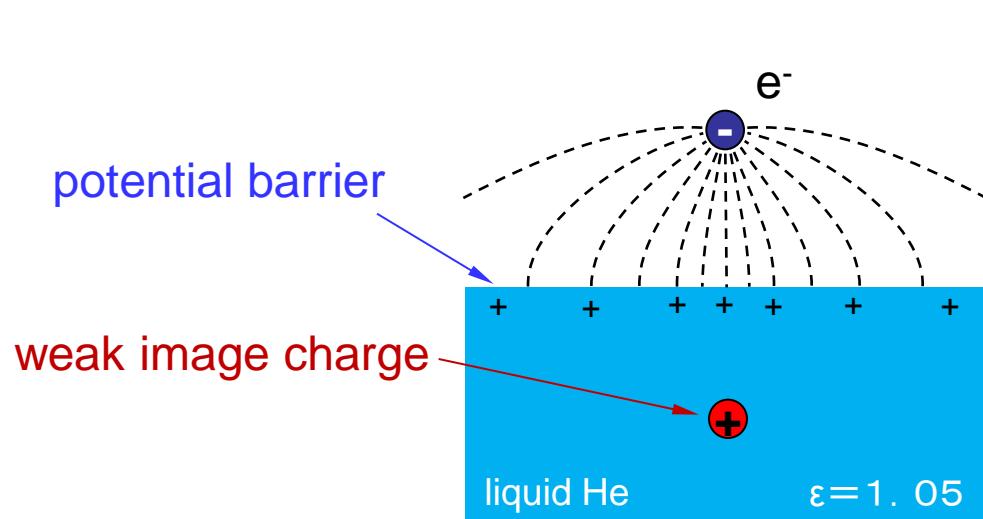


# Electrons on helium under CR excitation



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Transverse confinement

Image potential + surface barrier

Lateral confinement

Electrostatic potential due to  
positively biased gate

$$\Gamma_p = \frac{V_C}{E_{th}} = \frac{e^2 \sqrt{\pi n_s}}{4\pi\epsilon\epsilon_0 k_B T}$$

*Thermal fluctuations*

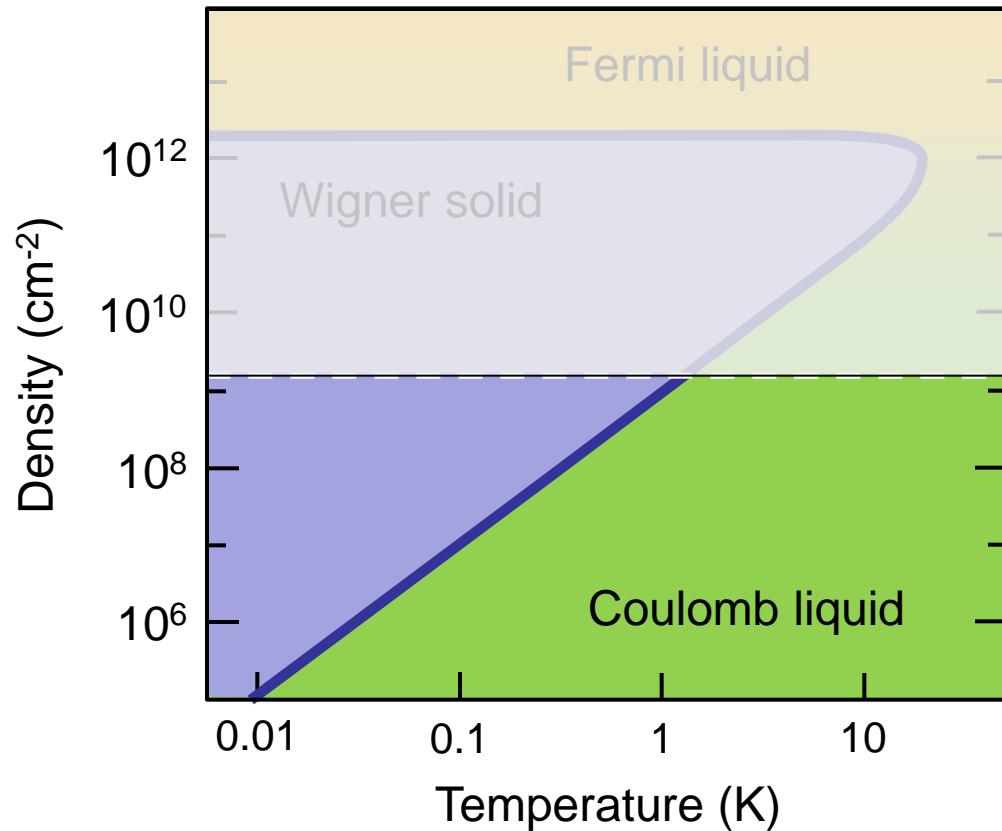
$$\epsilon=1$$

$$r_p = \frac{E_f}{V_C} = \frac{4\pi\epsilon\epsilon_0}{e^2 \sqrt{\pi n_s}} \frac{\hbar^2 \pi n_s}{m^*}$$

*Quantum fluctuations*

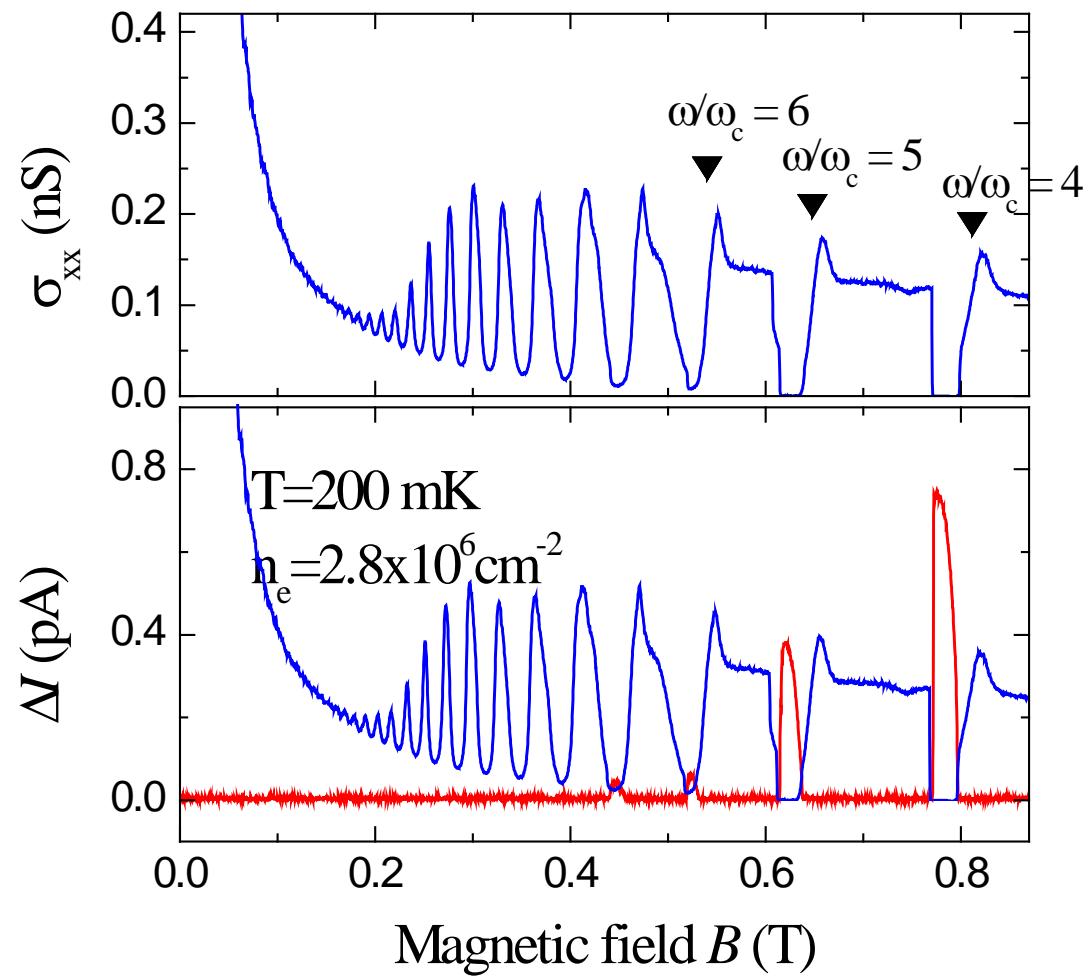
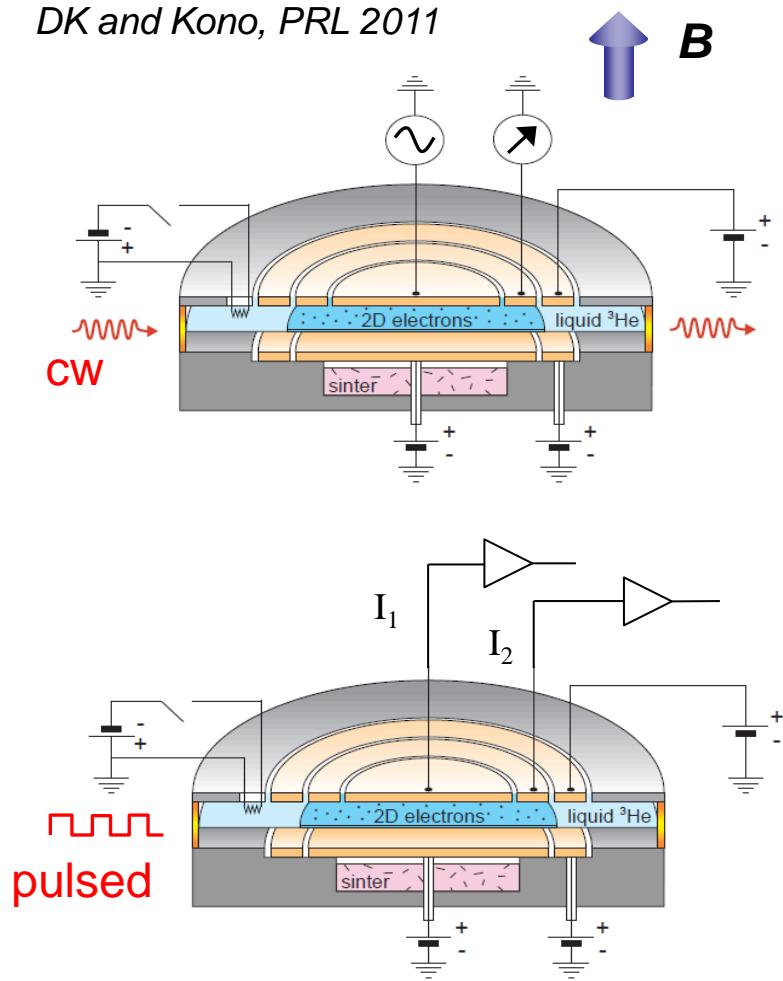
$$m^* = m_e$$

$$n_s < 2 \times 10^9 \text{ cm}^{-2}$$



2DEG in  
semiconductors

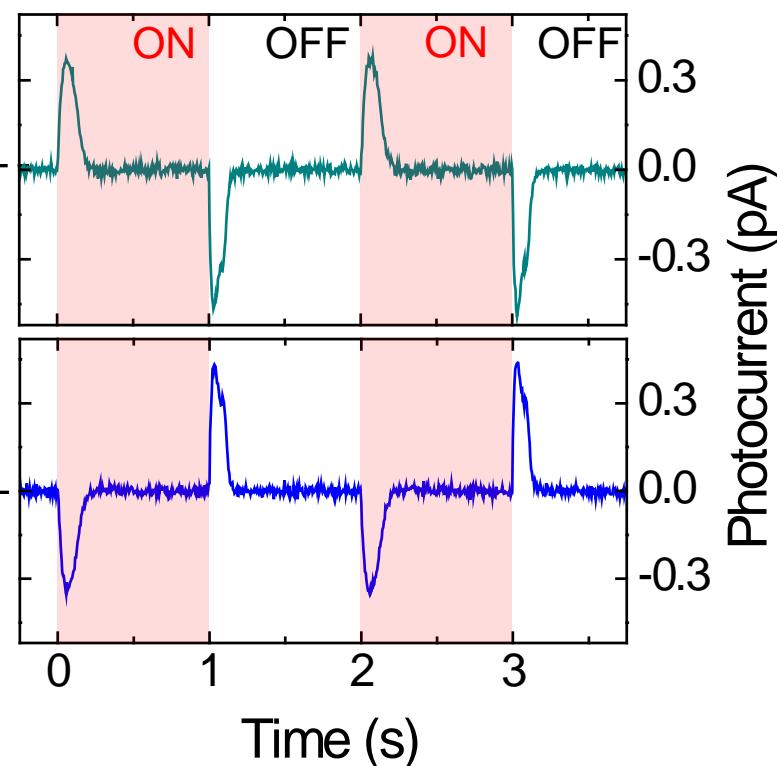
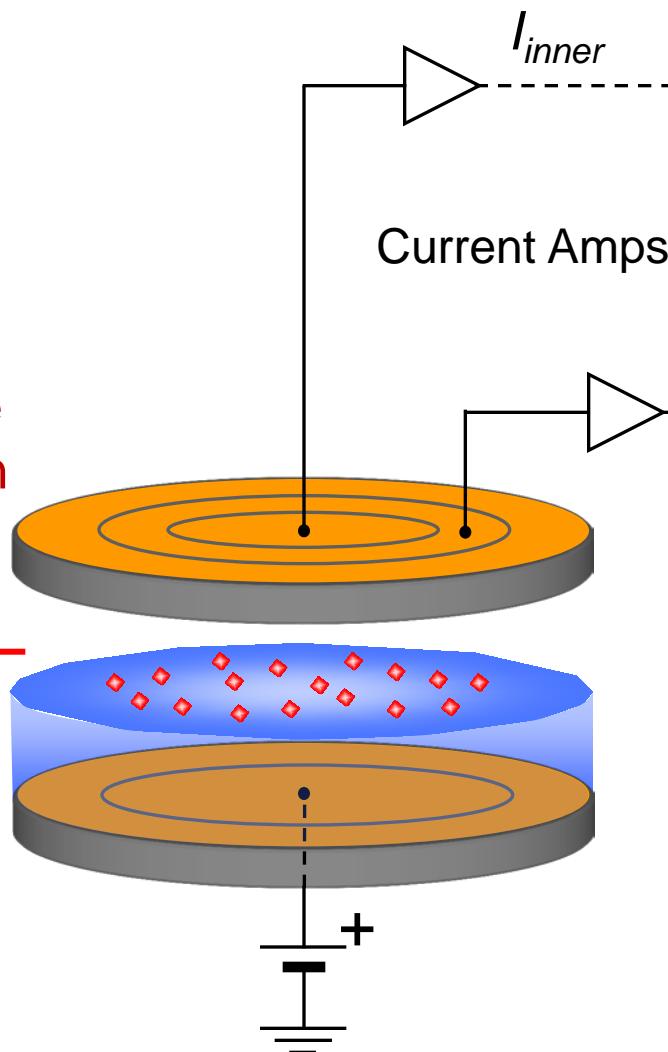
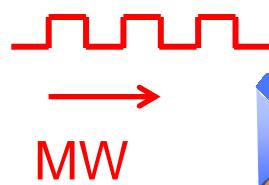
$\left. \begin{array}{l} \epsilon \approx 10, m^* \approx 0.1 m_e \\ n_s > 5 \times 10^{10} \text{ cm}^{-2} \end{array} \right\}$

*DK and Kono, PRL 2011*


DK, Chepelianskii, Kono, J. Phys. Soc. Jpn. 2012

Fix  $B$ -field  
at minima  
of  $\sigma_{xx}$

Microwave  
modulation  
ON/OFF



When MW switched ON  
Positive image charge released  
Electrons move away from the center

Momentum balance equation

Yu. P. Monarkha

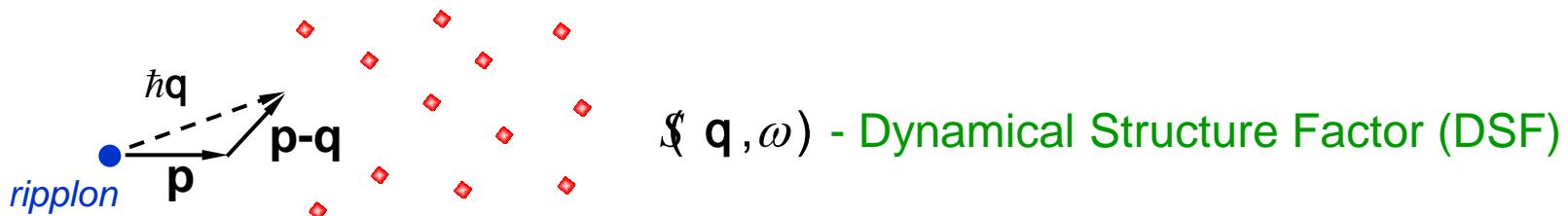


$$\mathbf{F}_{friction} = -N_e m v_{eff} \mathbf{u}_d$$

$$\sigma_{xx} = \frac{e^2 n_s}{m_e} \frac{v_{eff}}{\omega_c^2 + v_{eff}^2}$$

$$\mathbf{F}_{friction} = -\sum_{\mathbf{q}} \hbar \mathbf{q} \cdot P_{\mathbf{p},\mathbf{p}-\mathbf{q}} = -\frac{N_e}{\hbar^2 S} \sum_{\mathbf{q}} \hbar \mathbf{q} \cdot U(\mathbf{q}) \cdot \underbrace{S(\mathbf{q}, \omega - \mathbf{q} \cdot \mathbf{u}_d)}$$

Doppler shift in laboratory frame



$$\mathbf{F}_{friction} = -\frac{N_e}{\hbar^2 S} \sum_{\mathbf{q}} \hbar \mathbf{q} \sum_{n,n'} \rho_n U_{n,n'}(\mathbf{q}) \left[ \underbrace{(N_q + 1) S_{n,n'}(\mathbf{q}, \omega_{n,n'} - \mathbf{q} \cdot \mathbf{u}_d)}_{\text{creation of ripplon}} - \underbrace{N_q S_{n,n'}(\mathbf{q}, \omega_{n,n'} + \mathbf{q} \cdot \mathbf{u}_d)}_{\text{annihilation of ripplon}} \right]$$

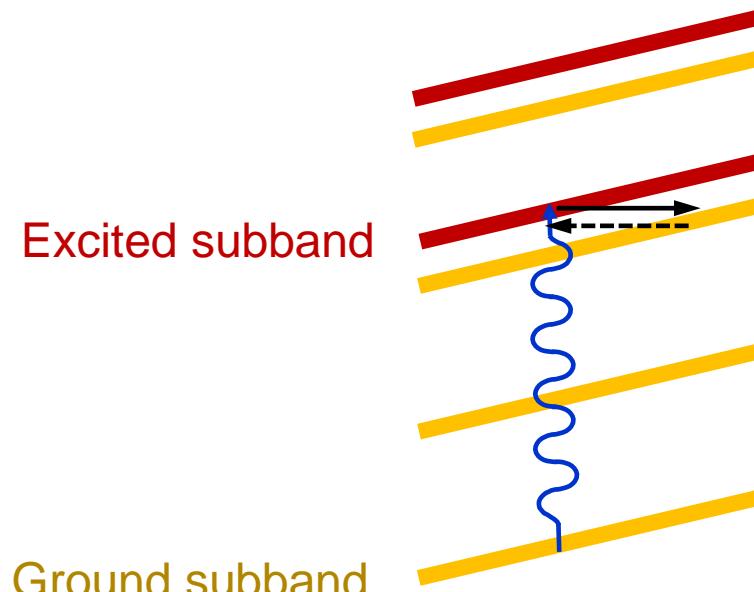
$$\mathbf{F}_{friction} \sim -\sum_{\mathbf{q}} \hbar \mathbf{q} \sum_{n,n'} \chi_{n,n'}(\mathbf{q}) \cdot S_{n,n'}(\mathbf{q}, \omega_{n,n'} + \mathbf{q} \cdot \mathbf{u}_d) \cdot (\rho_n - \rho_{n'} e^{-\frac{\hbar \omega_{n,n'}}{k_B T_e}} e^{-\frac{\hbar \mathbf{q} \cdot \mathbf{u}_d}{k_B T_e}})$$

Yu. P. Monarkha

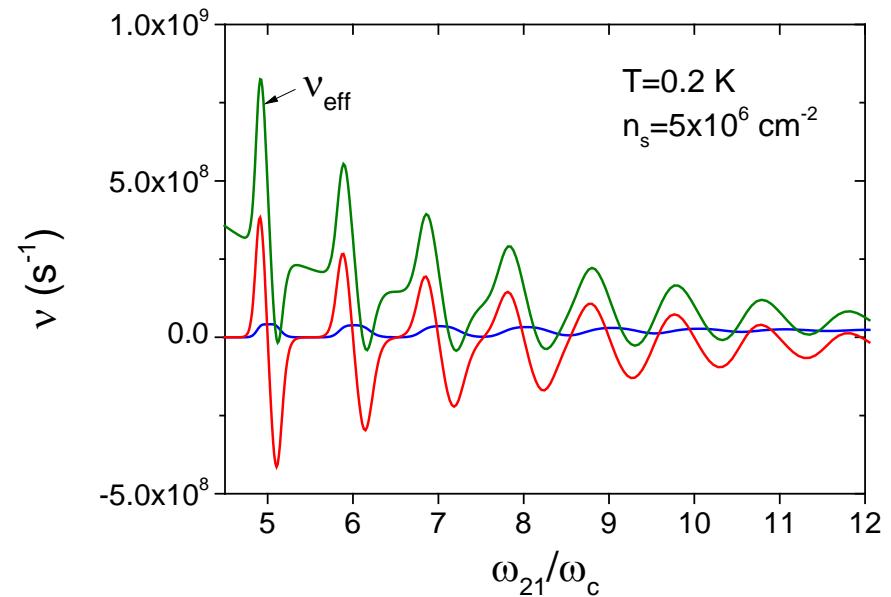
$$\nu_{\text{inter}} \sim \sum_{\mathbf{q}} \hbar \mathbf{q} \sum_{n>n'} \chi_{n,n'}(\mathbf{q}) \cdot \underbrace{S_{n,n'}(q, \omega_{n,n'})}_{\text{DSF}} \cdot (\rho_n + \rho_{n'} e^{-\frac{\hbar \omega_{n,n'}}{k_B T_e}}) +$$

$$+ \frac{\hbar}{k_B T} \sum_{\mathbf{q}} \hbar \mathbf{q} \sum_{n>n'} \chi_{n,n'}(\mathbf{q}) \cdot \underbrace{S'_{n,n'}(q, \omega_{n,n'})}_{\text{derivative of DSF}} \cdot (\rho_n - \rho_{n'} e^{-\frac{\hbar \omega_{n,n'}}{k_B T_e}})$$

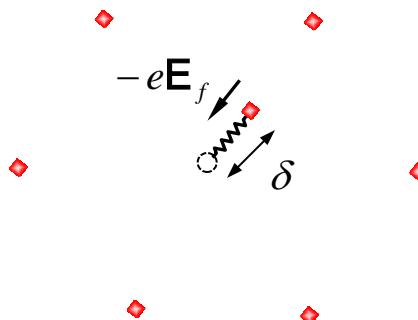
Tilted LLs in dc E-field



non-zero for nonequilibrium population



Dykman, Khazan 1979

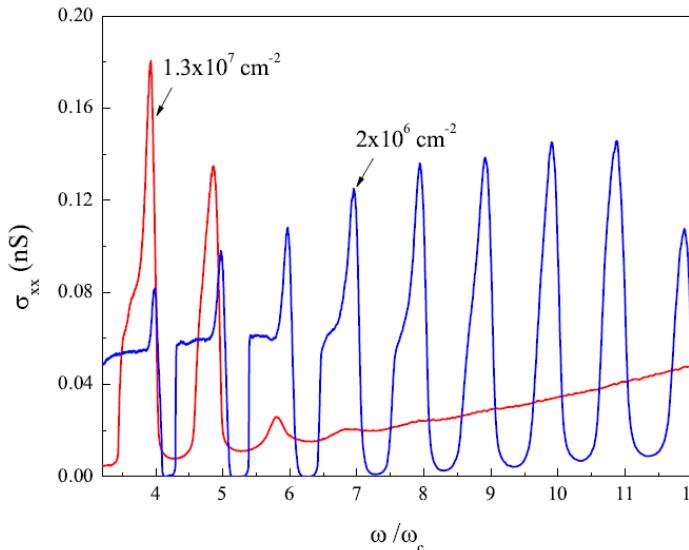


$$\frac{k\delta^2}{2} \approx k_B T_e \quad \Rightarrow \quad \langle E_f \rangle \approx n_s^{3/4} \sqrt{4\pi k_B T_e}$$

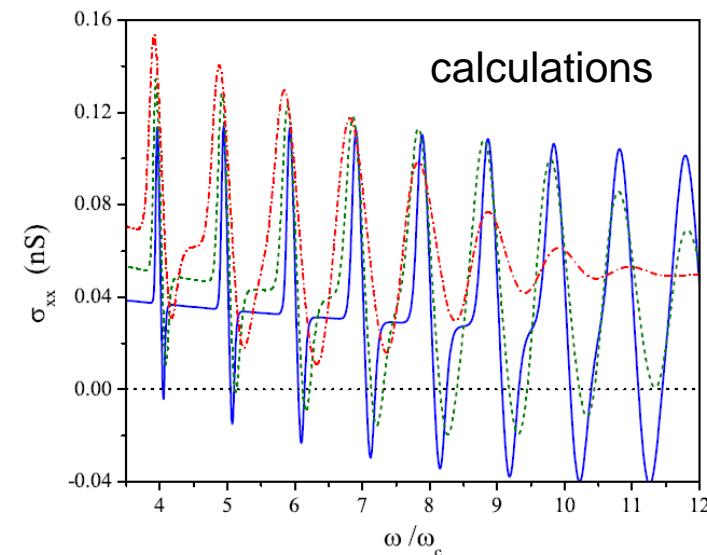
$$\mathbf{u}_d \rightarrow \mathbf{u}_f$$

$$S^{(me)}(q, \omega) = \langle S(q, \omega - \mathbf{q} \cdot \mathbf{u}_f) \rangle$$

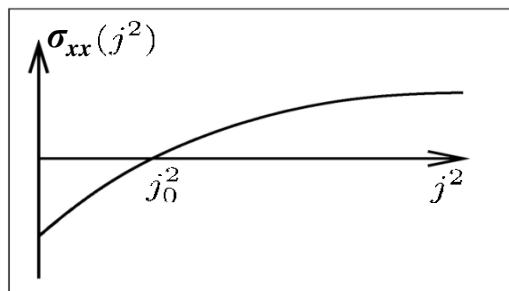
DK, Monarkha, Kono, PRL 2013



Coulomb broadening of DSF



There exists (some) mechanism  
leading to **absolute negative  $\sigma_{xx}$**



$$\frac{\partial \rho_s}{\partial t} = -\nabla \vec{j} = \sigma \nabla_r^2 V$$

$$\nabla^2 V = -\rho_s \delta(z)$$

$$\rho_k \sim \exp\left(-\frac{\sigma k}{2} t\right) \rightarrow \infty$$

**Instability and formation of current domains**

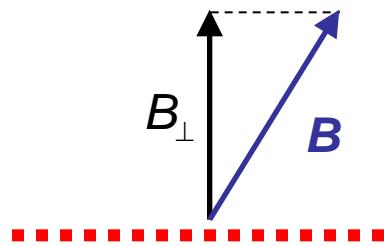
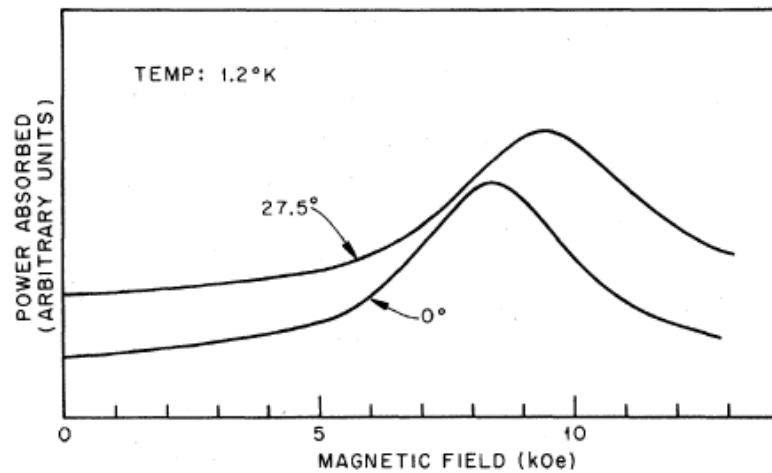
*Andreev, Aleiner, Millis, PRL 2003*

Other explanations are proposed, e.g. Entin and Magarill: [arXiv: 1504.03422](https://arxiv.org/abs/1504.03422)

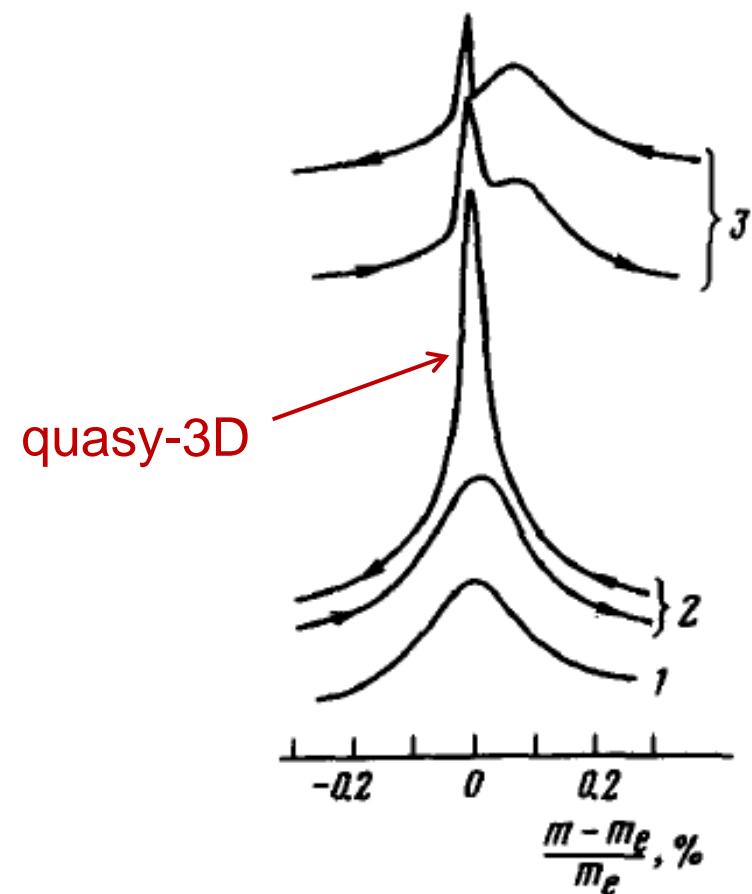
ZRS in electrons on helium have interesting properties

- self-organized oscillations (previous talk)
- incompressibility (next talk)

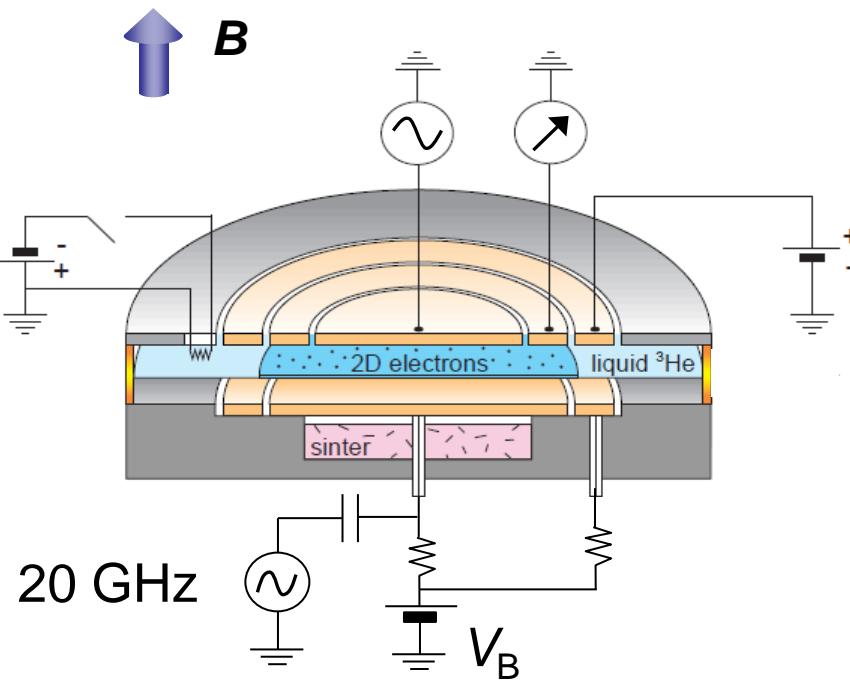
Brown, Grimes 1972



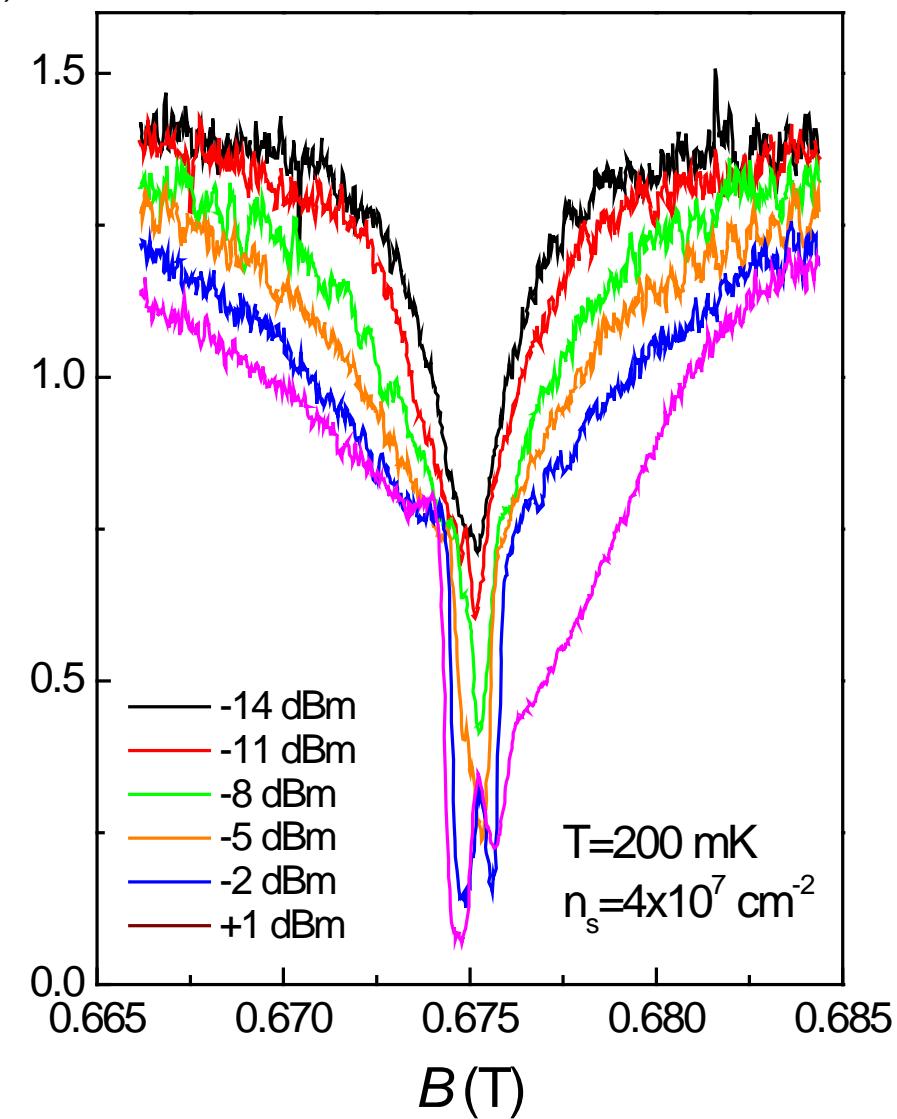
Edelman 1972



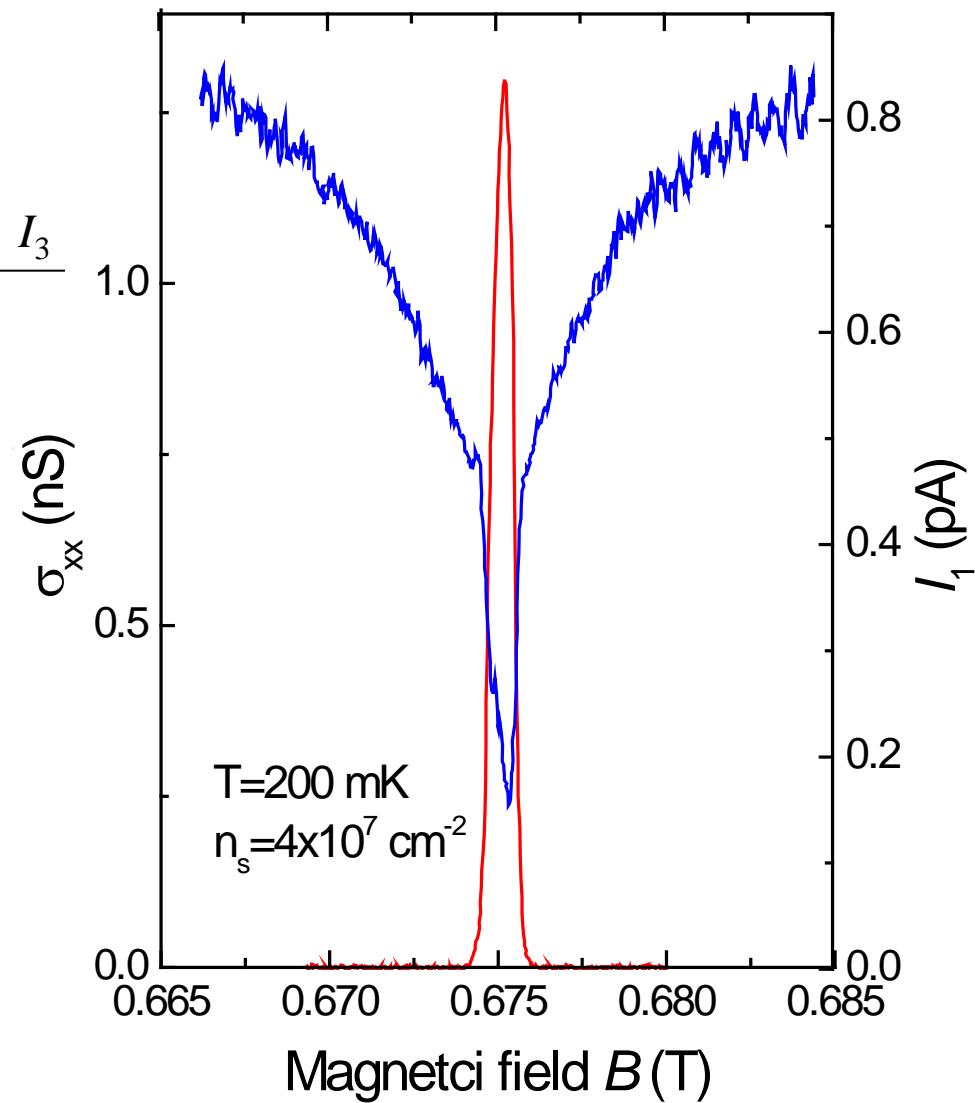
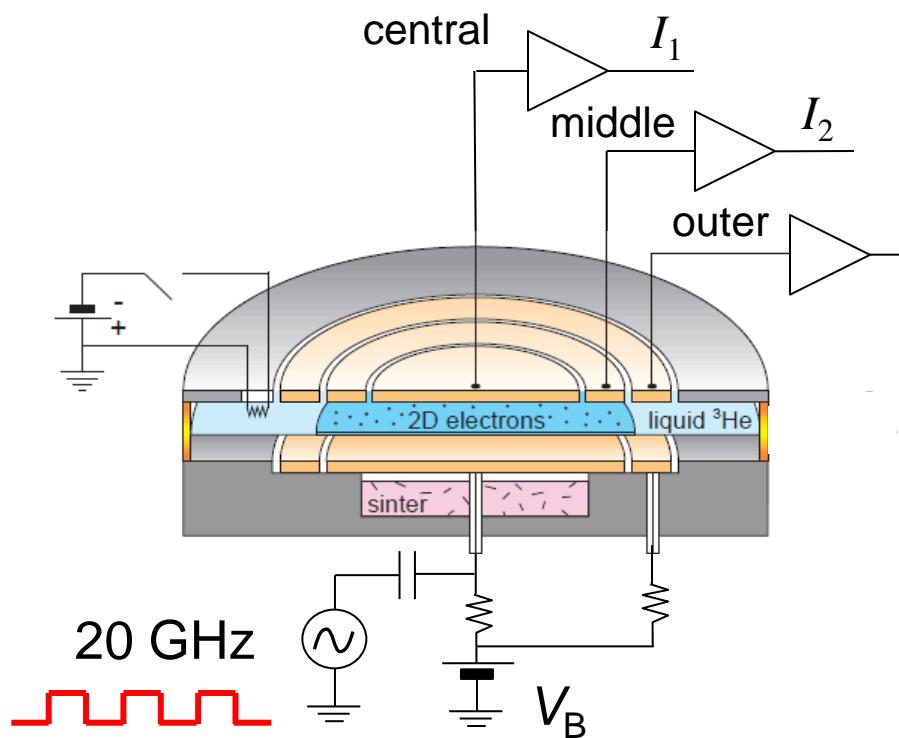
Badrutdinov , Abdurakhimov, DK, PRB 90, 075305 (2014)

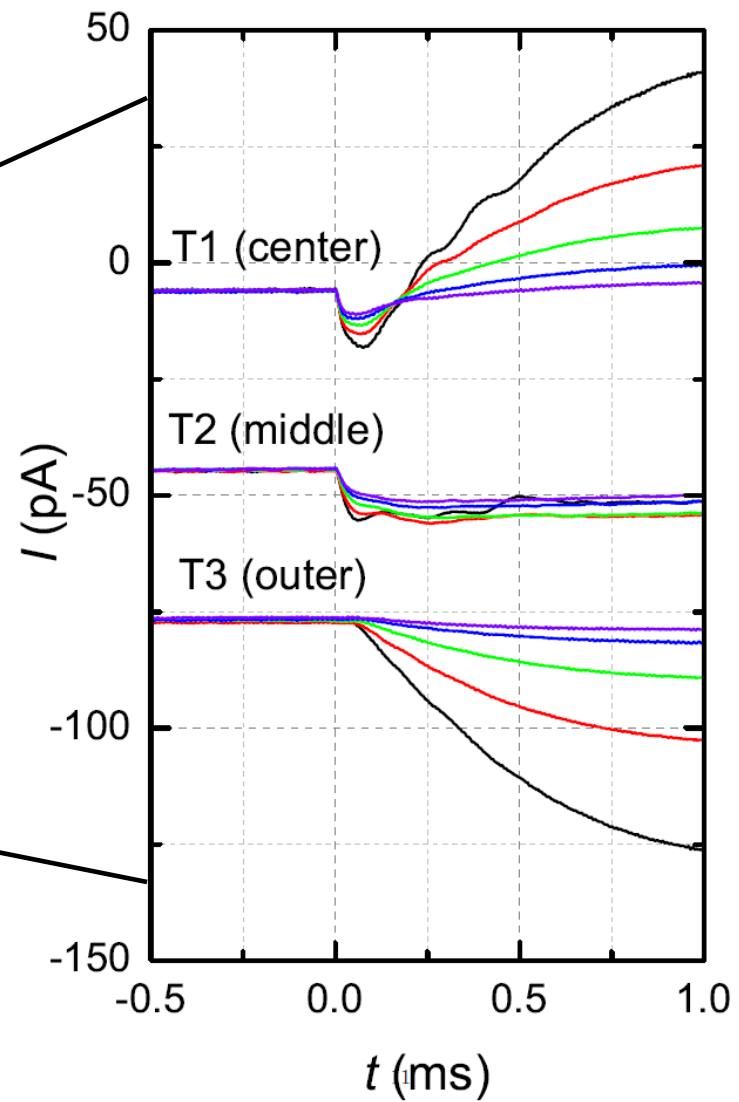
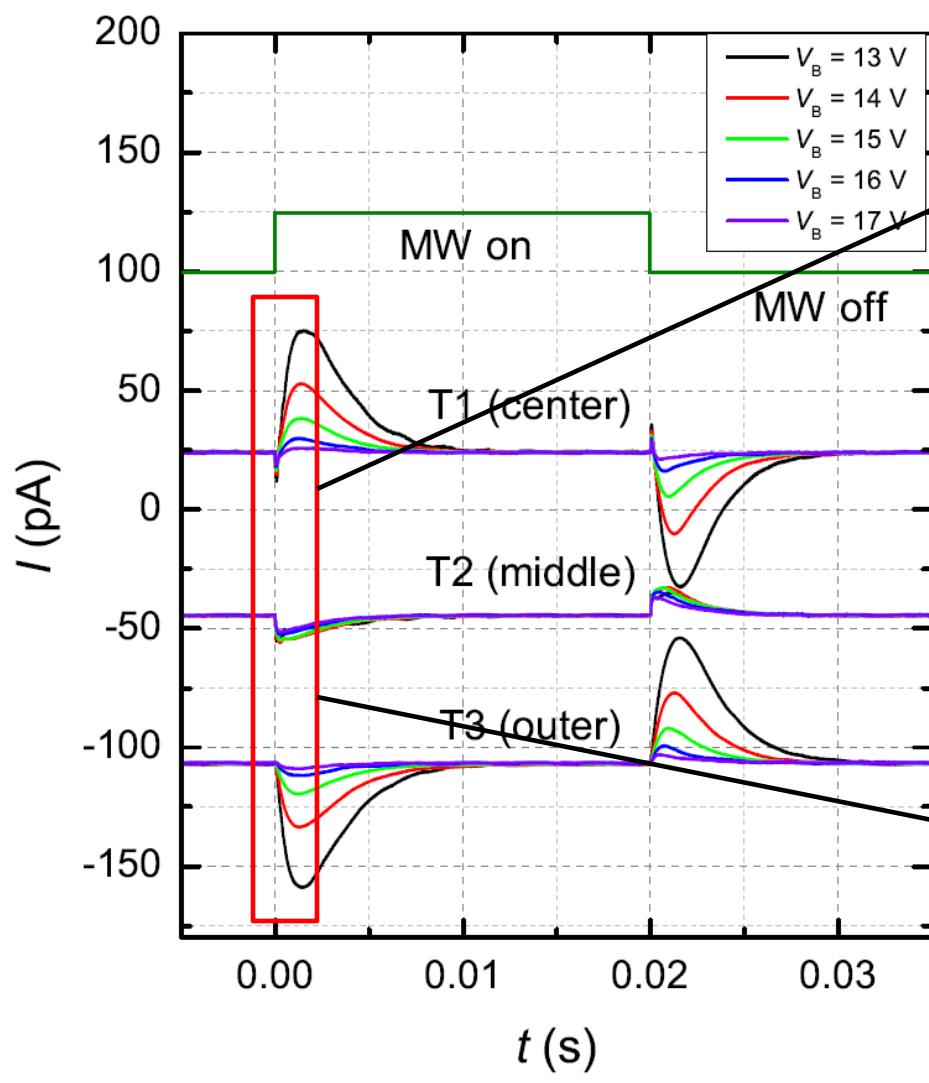


Irreversible loss of electrons  
at high MW power



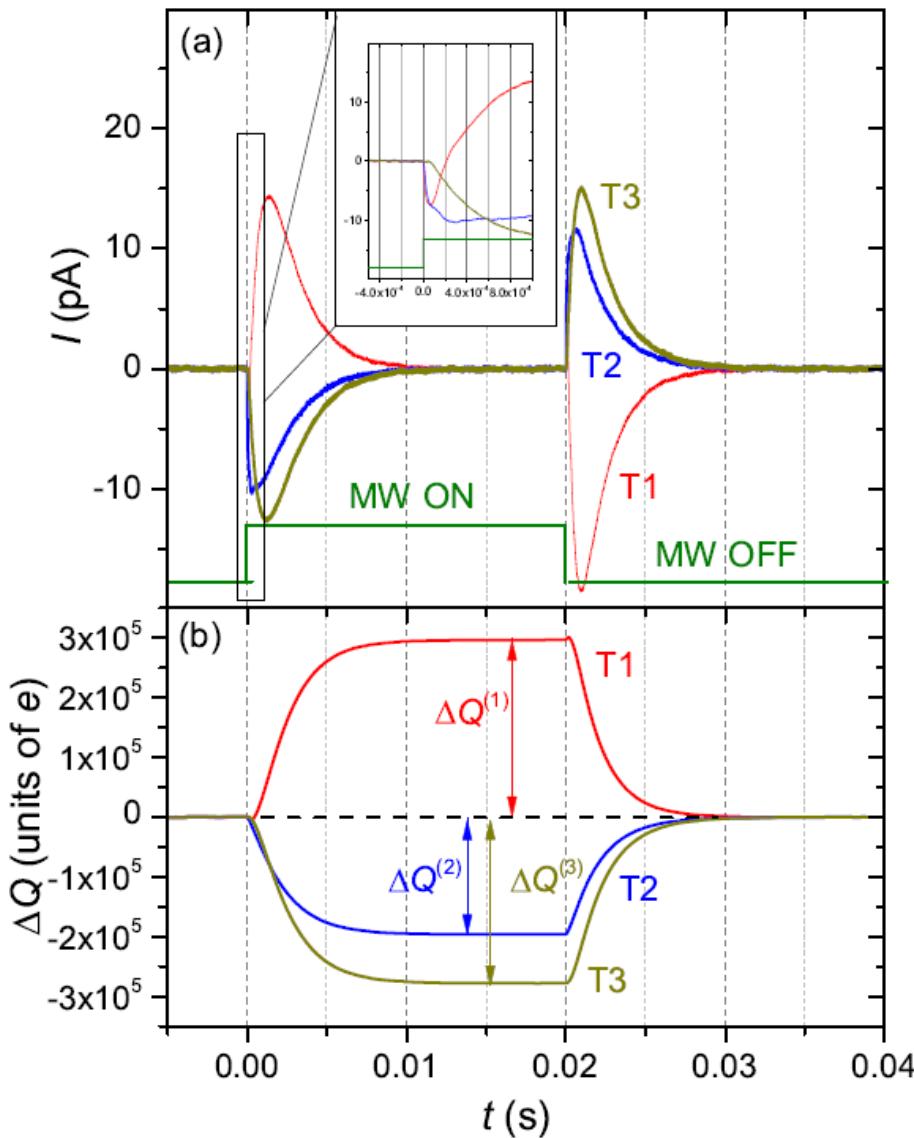
Badrutdinov , Abdurakhimov, DK, PRB 90, 075305 (2014)







Badrutdinov , Abdurakhimov, DK, PRB 90, 075305 (2014)



Extraction of  $\Delta Q_{\text{perp}}^{(i)}$  and  $\Delta Q_{\text{in-plane}}^{(i)}$  from data:

$$\Delta Q^{(i)} = \Delta Q_{\text{in-plane}}^{(i)} + \Delta Q_{\text{perp}}^{(i)}$$

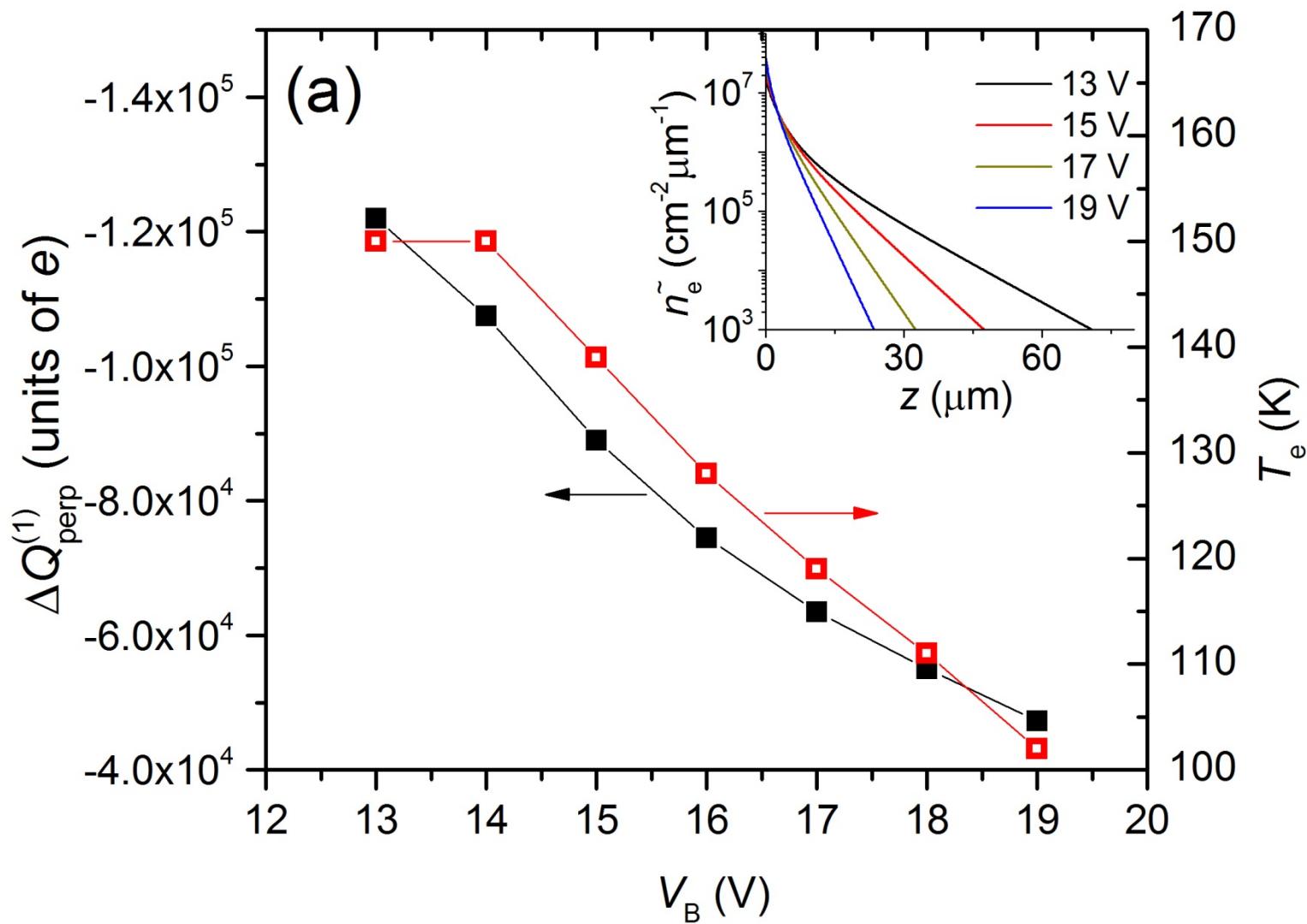
$$\Delta Q_{\text{in-plane}}^{(1)} + \Delta Q_{\text{in-plane}}^{(2)} + \Delta Q_{\text{in-plane}}^{(3)} = 0$$

$$\Delta Q_{\text{perp}}^{(1)} = \Delta Q_{\text{perp}}^{(2)}$$

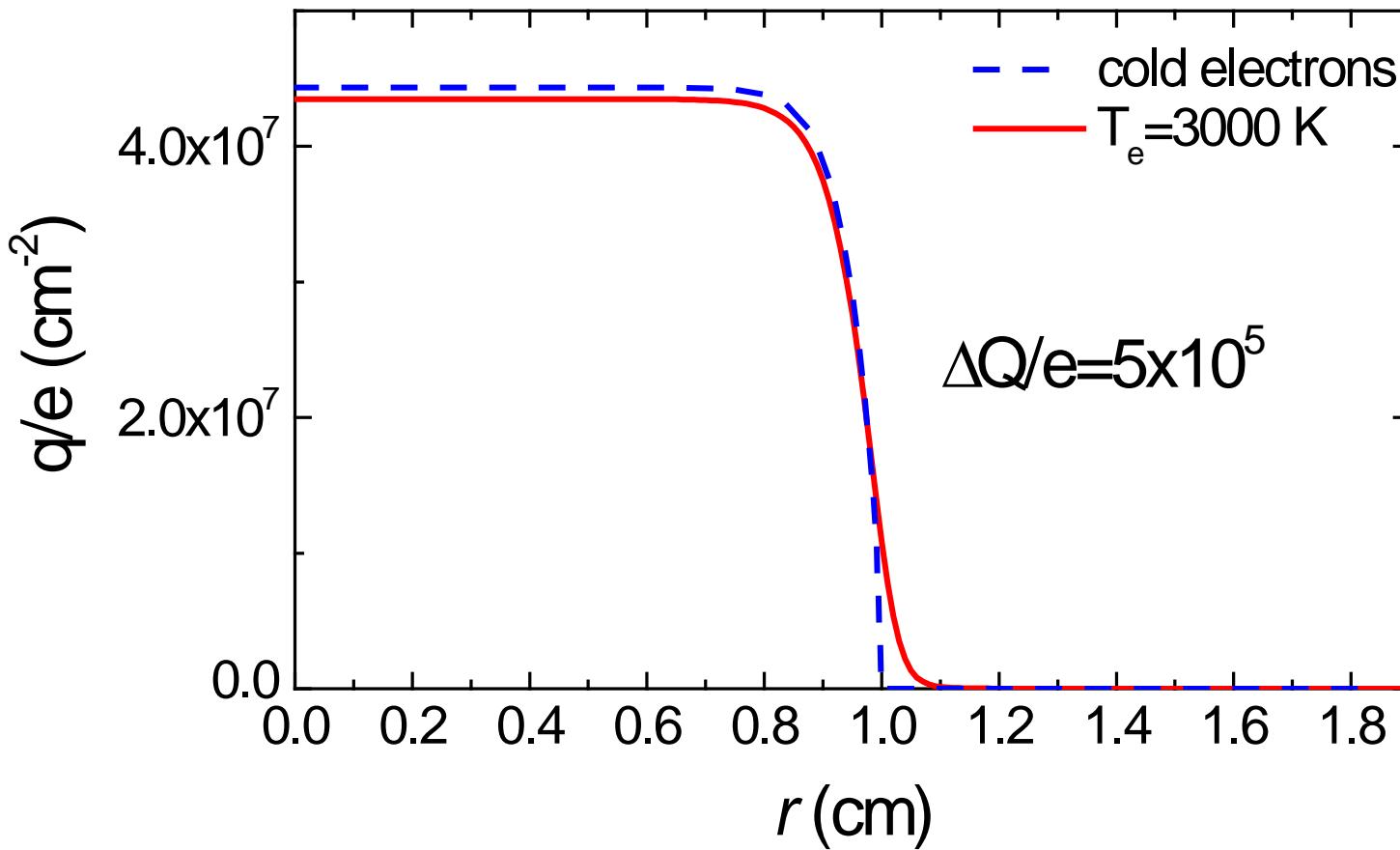
$$\Delta Q_{\text{perp}}^{(3)} = 0$$

$V_B$ (v)	$\Delta Q_{\text{in-plane}}^{(1)}$	$\Delta Q_{\text{in-plane}}^{(2)}$	$\Delta Q_{\text{in-plane}}^{(3)}$	$\Delta Q_{\text{perp}}$
13	-1338000	55000	1283000	122000
14	-736500	109500	627000	107500
15	-385000	107000	278000	89000
16	-184500	72500	112000	74500
17	-95500	46500	49000	63500
18	-56800	32000	24800	55000
19	-38600	23700	14900	47300

$$Q_0 = 1.35 \cdot 10^8$$

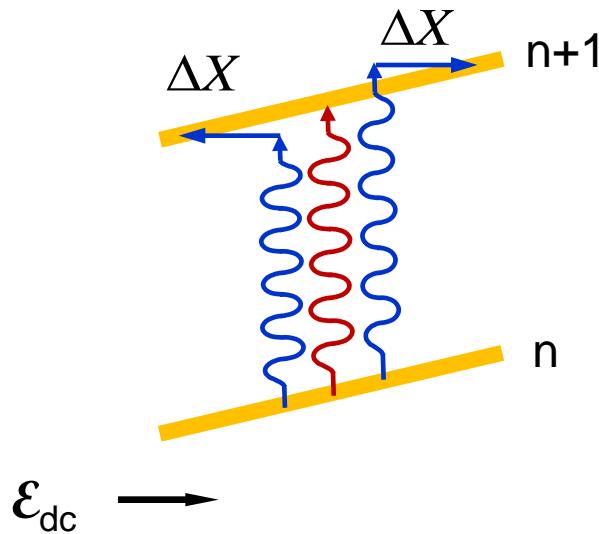


Simultaneous solution of Poisson and Boltzmann equation gives estimate of 3,000 K for electron temperature



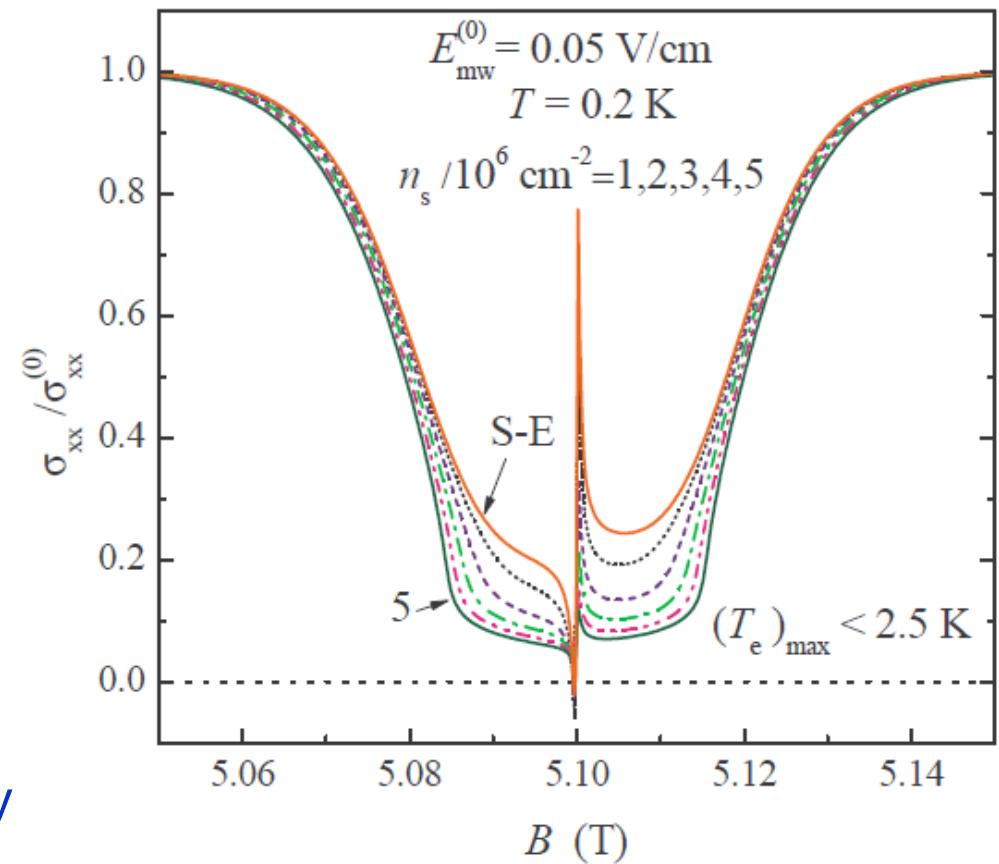
## Breakdown of the effective electron temperature approximation?

Yu. Monarkha, PRB(R) 2015

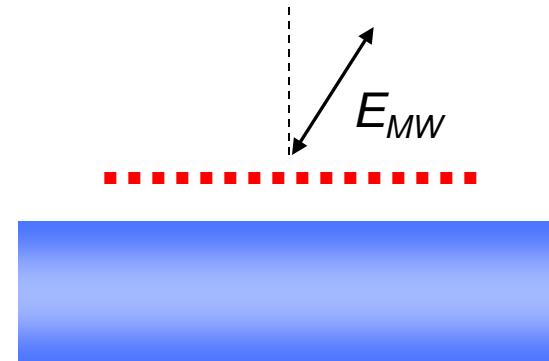


Predicts negative  $\sigma_{xx}$  and instability

on the low-field side of CR

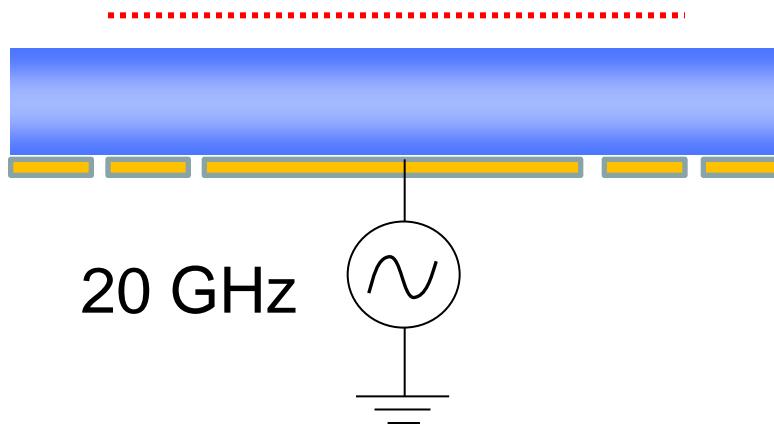


*Magarill and Entin, arXive:1504.03422*



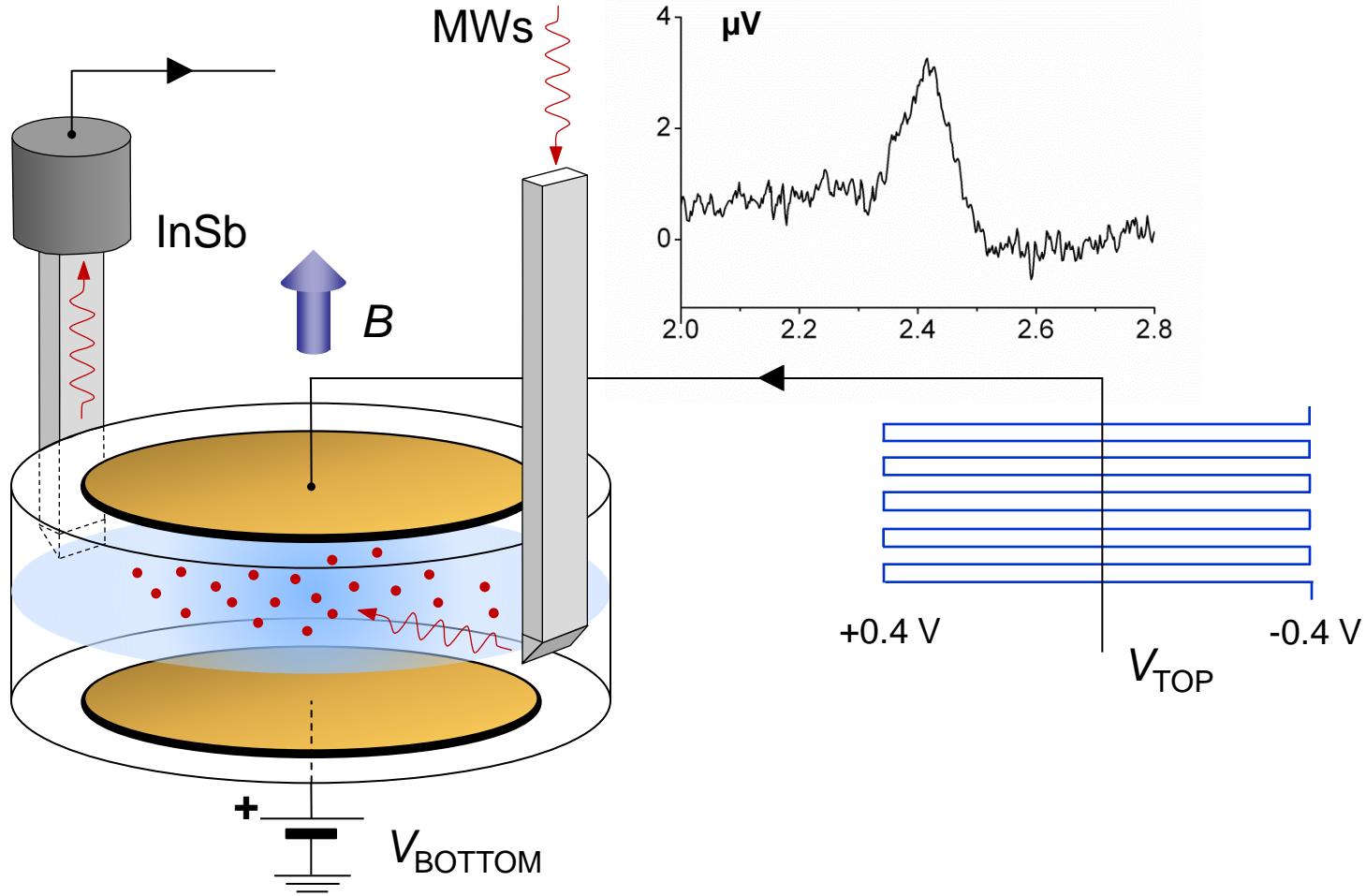
$$j_x = \left\langle e n_s \sum_{\beta'} (X_{\beta'} - X_{\beta}) w_{\beta \rightarrow \beta'} \right\rangle$$

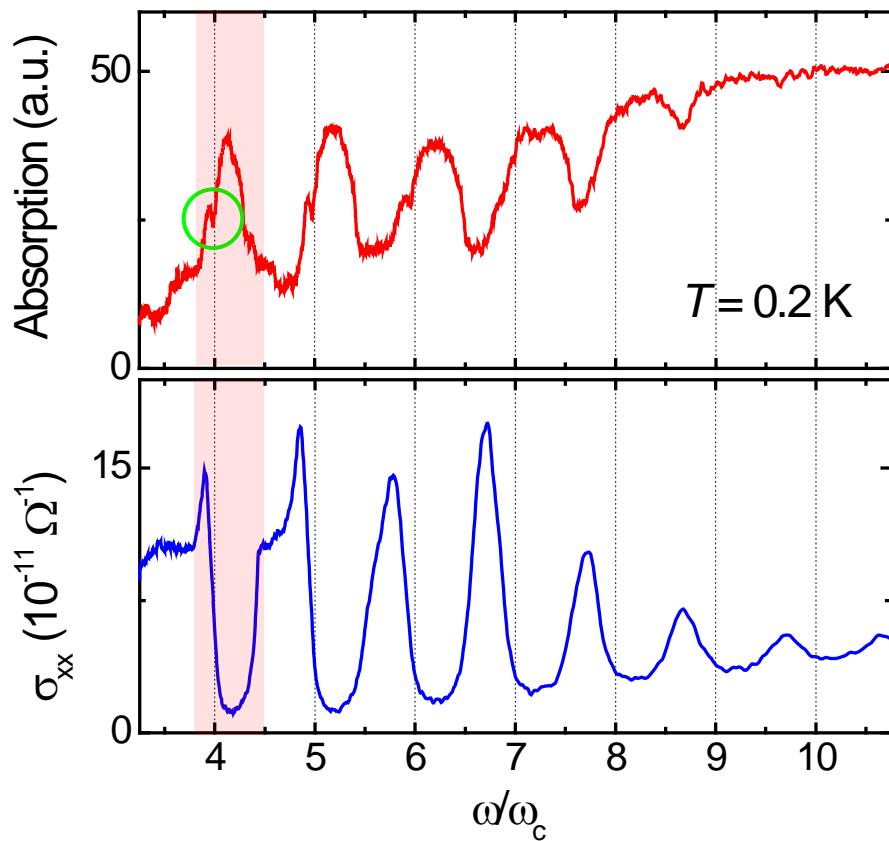
Second order perturbation theory



Would expect radial photocurrent  
in our geometry

- In principle, our observation is consistent with formation of a quasy-3D electron system speculated earlier.
- However, it is difficult to understand why different motional degrees of freedom do not thermalize.
- Other theoretical proposals might be relevant but require examination





Energy balance equation:

$$\text{Absorption} = v_E \cdot (T_e - T)$$

Heating of electrons:

