

Realization of spin qubits using electrons on the surface of superfluid helium

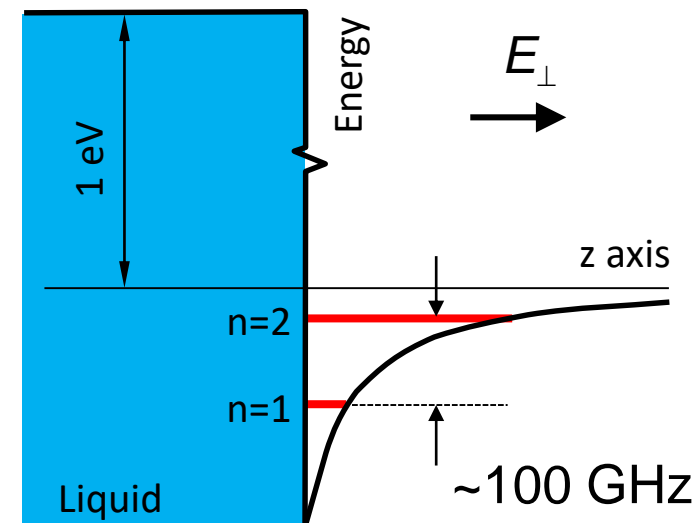
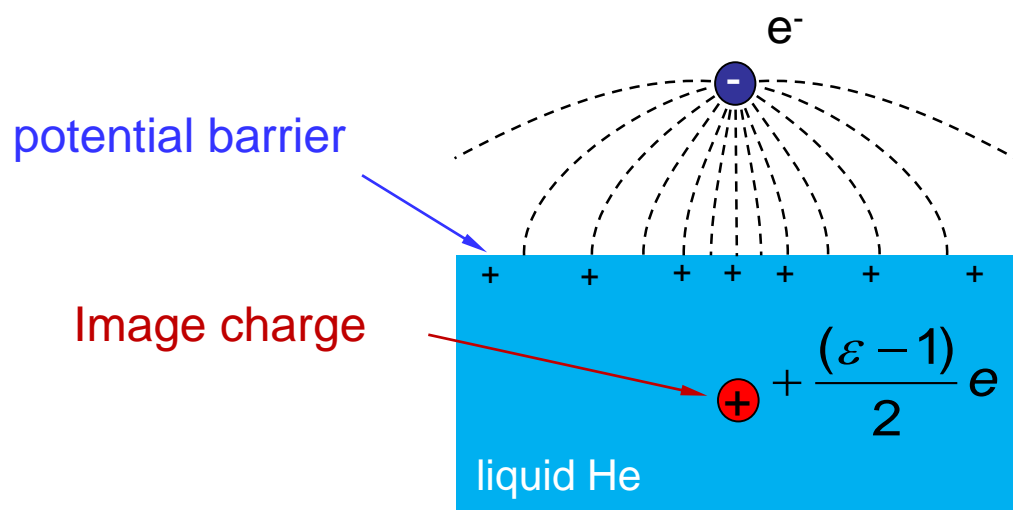
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OKINAWA INSTITUTE OF SCIENCE AND TECHNOLOGY GRADUATE UNIVERSITY

APTQS 2022

Surface States of Electrons on Helium (EonHe)



Hydrogen-like (Rydberg) spectrum:

$$E_n = -\frac{R_y}{n^2}, \quad n = 1, 2, \dots$$

Small Rydberg energy:

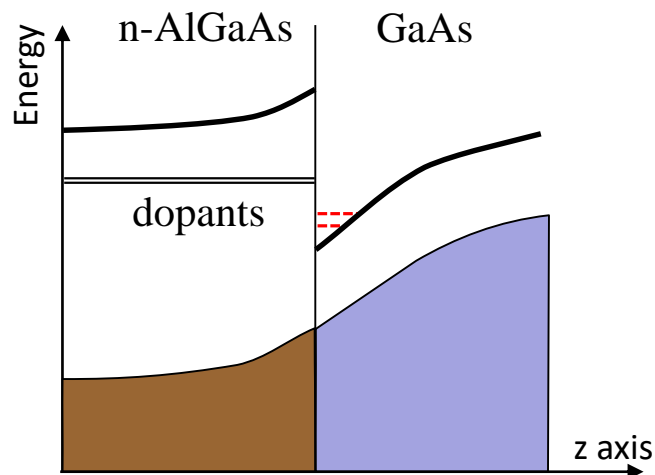
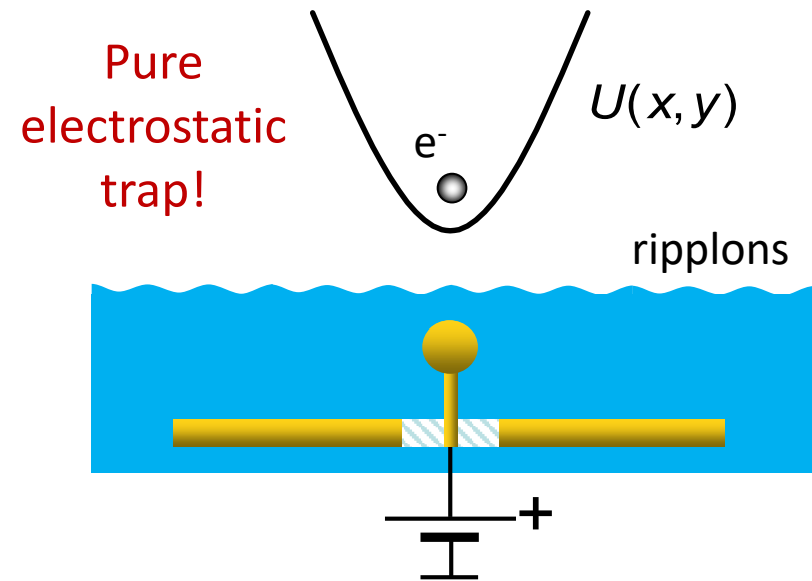
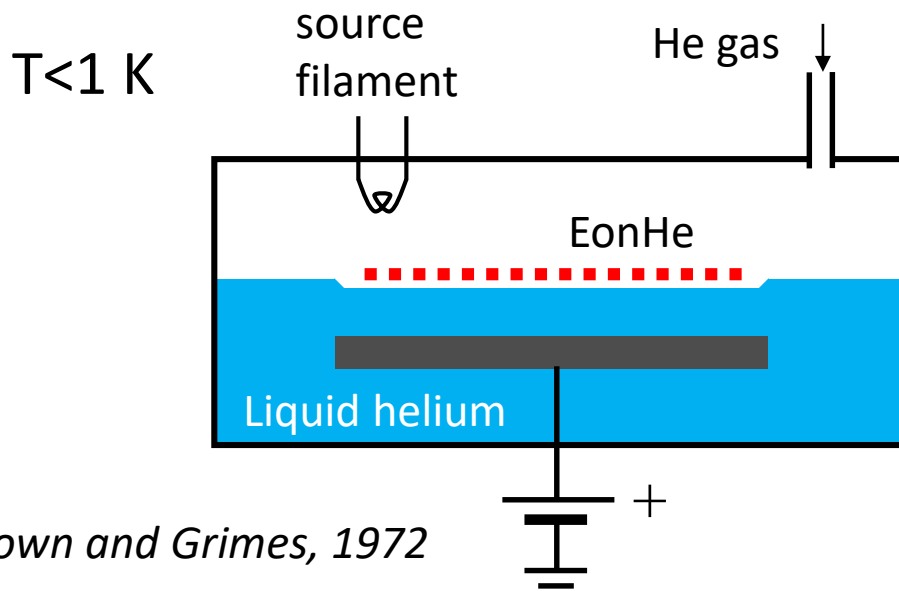
$$R_y = \frac{(\epsilon - 1)^2 m_e^2 e^2}{16\hbar^2} \approx 10^{-3} \text{ eV} = 10 \text{ K}$$

Linear Stark shift:

$$\Delta E_n = eE_{\perp} z_{nn}$$



2D electron system



Why liquid helium?

- Remains liquid down to $T=0$
- Smooth surface, no defects
- Interaction only with [ripples](#)



Quantum computing using Rydberg states

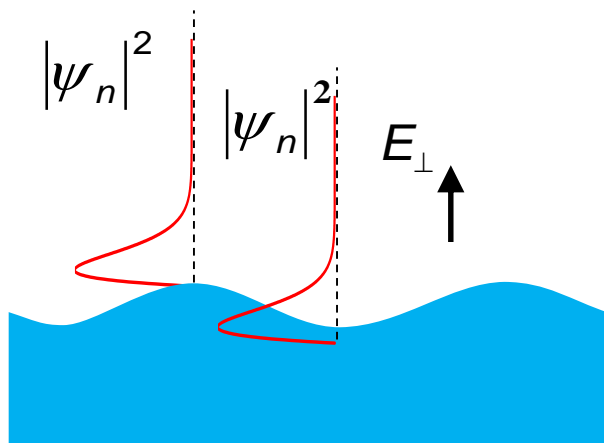
SCIENCE VOL 284 18 JUNE 1999

Quantum Computing with Electrons Floating on Liquid Helium

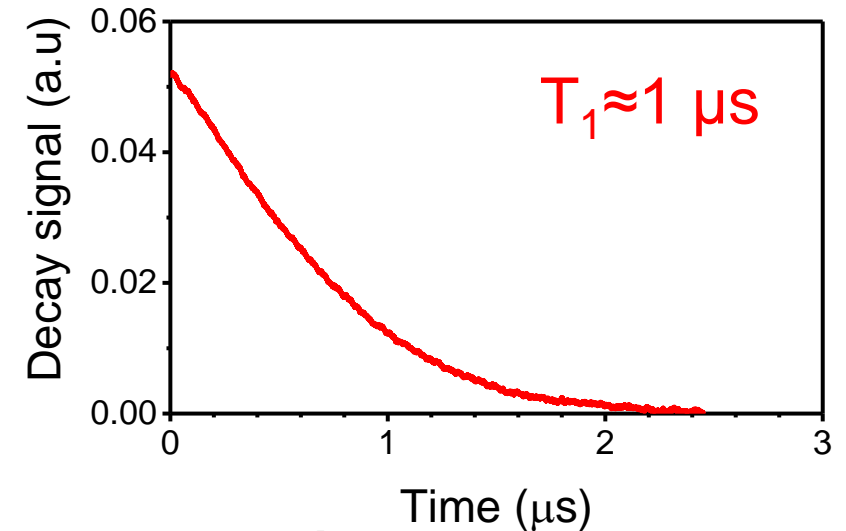
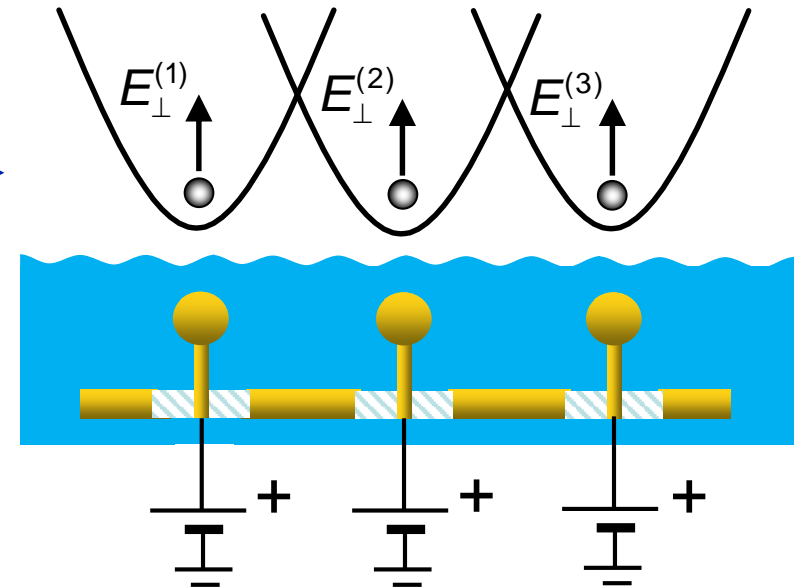
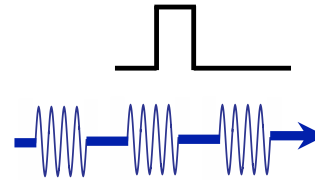
P. M. Platzman^{1*} and M. I. Dykman²

A quasi-two-dimensional set of electrons ($1 < N < 10^9$) in vacuum, trapped in one-dimensional hydrogenic levels above a micrometer-thick film of liquid helium, is proposed as an easily manipulated strongly interacting set of quantum bits. Individual electrons are laterally confined by micrometer-sized metal pads below the helium. Information is stored in the lowest hydrogenic levels. With electric fields, at temperatures of 10^{-2} kelvin, changes in the wave function can be made in nanoseconds. Wave function coherence times are 0.1 millisecond. The wave function is read out with an inverted dc voltage, which releases excited electrons from the surface.

~~$T_1 = 100 \mu\text{s}$~~



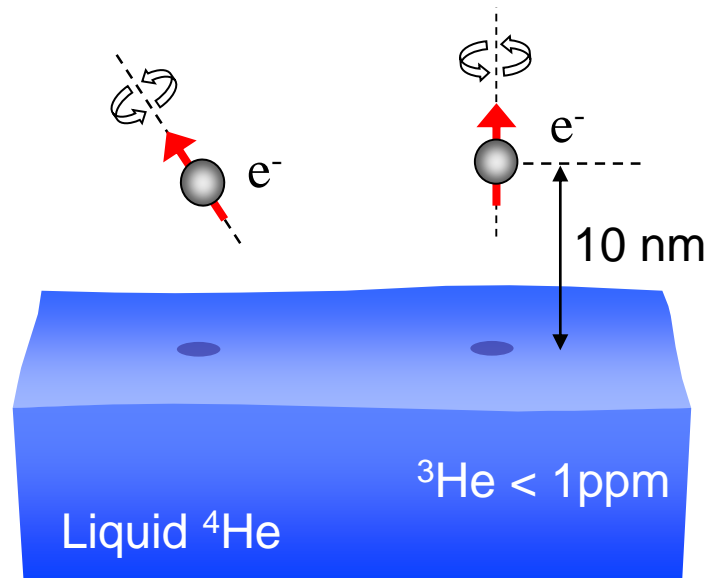
MW pulses



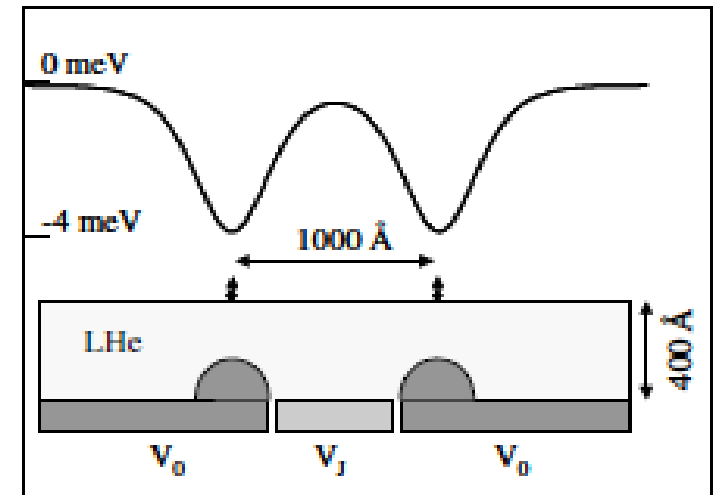
[E. Kawakami *et al.*, PRL 126, 106802]



Quantum computing using spins



Steve Lyon, 2004



[S. A. Lyon, Phys. Rev. A, **74**, 052338]

Spin coherence time $>100\text{ s!}$

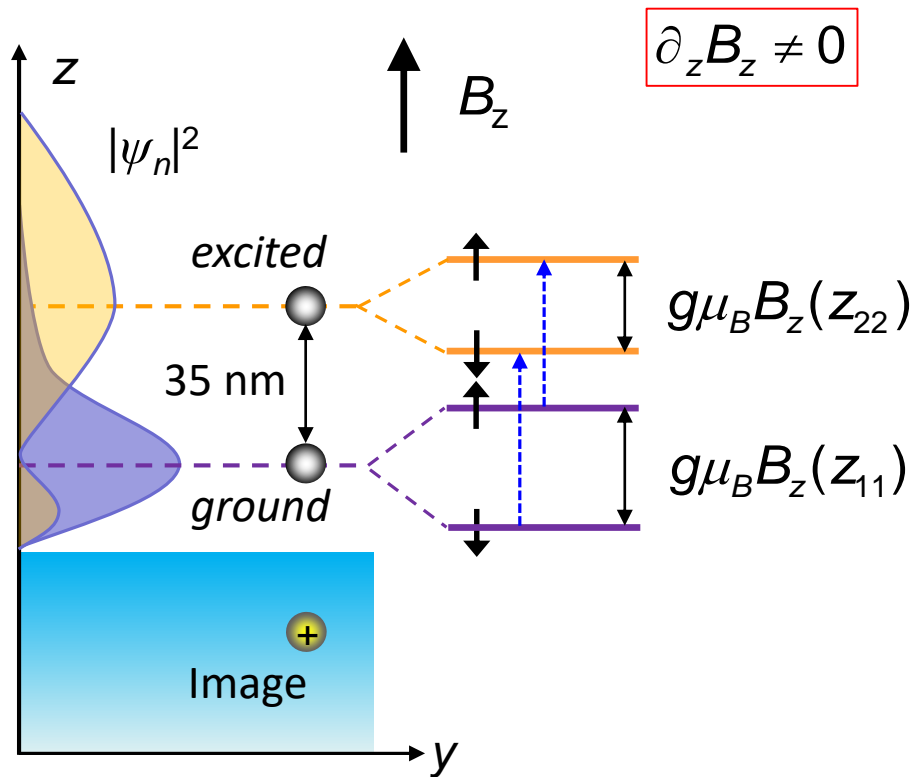
- Negligible SO interaction
- Magnetic-impurity-free environment
- Negligible noise from electrodes

How to control spin states?

- Very weak dipole interaction
- Slow spin rotations
- Spin state readout?



Introduce SO coupling (Spin-Rydberg)



Different Zeeman splitting for ground and excited Rydberg states!



Erika Kawakami

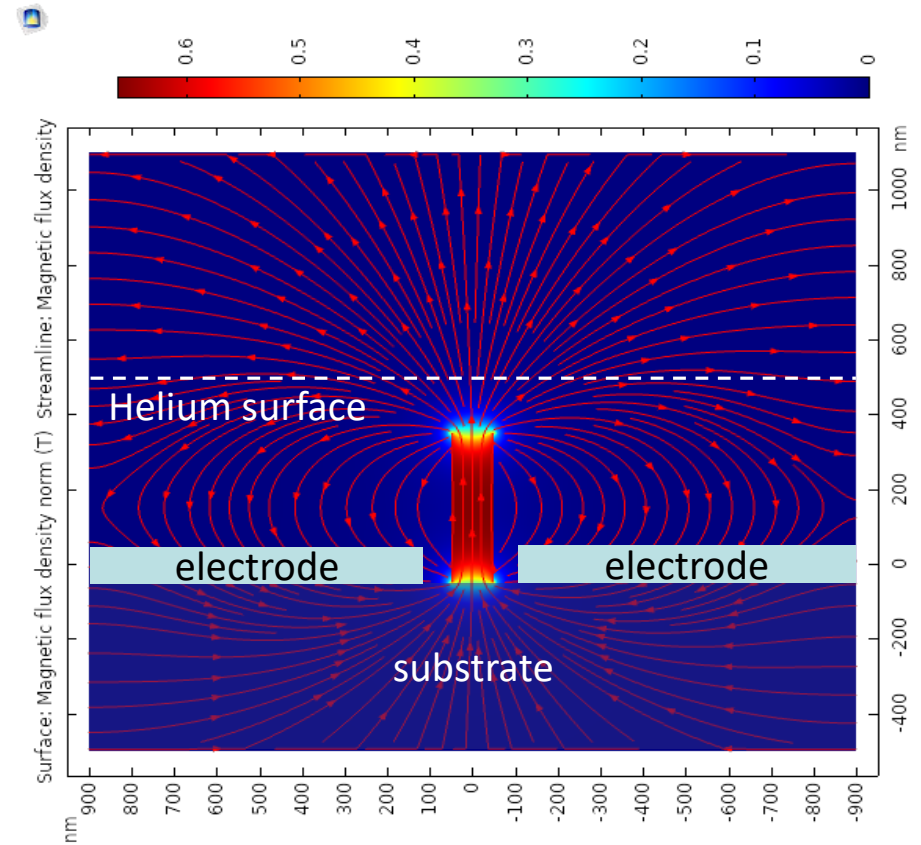
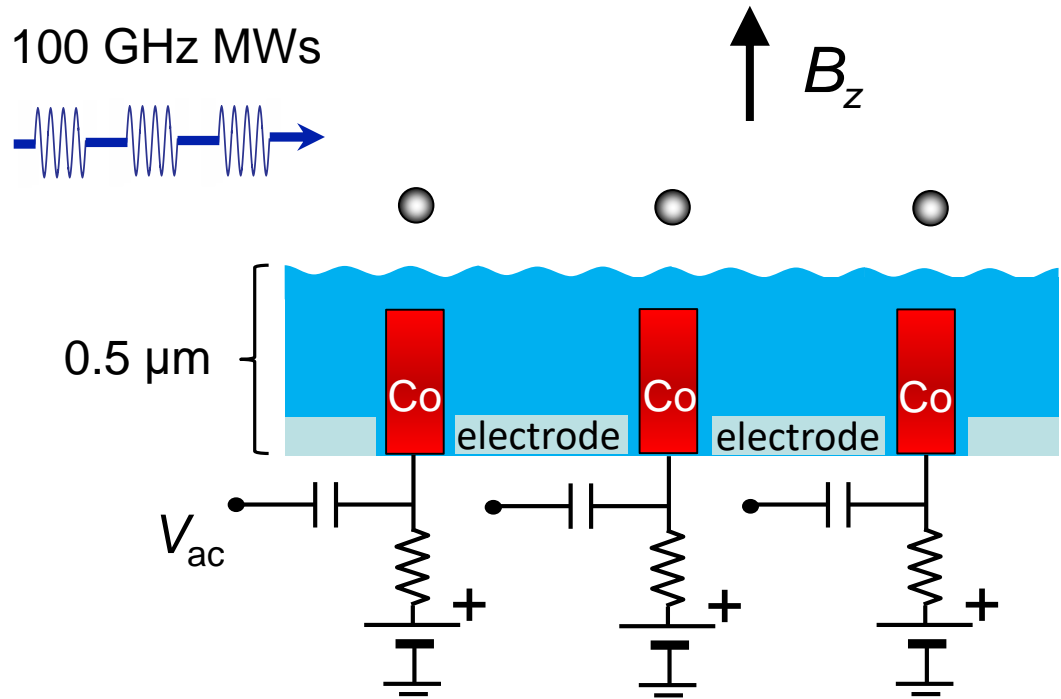
Spin-orbit (SO) interaction:

$$H_{SO} = \gamma_e \left(\frac{\partial B_z}{\partial z} \right) z s_z$$

- Fast spin rotations (single-q gate)
- Spin-state readout (q-readout)
- Spin-spin coupling (two-q gate)



EonHe-Spin quantum computer



Difference in Zeeman splitting:

$$g\mu_B \left(\frac{\partial B_z}{\partial z} \right) (z_{22} - z_{11}) \approx 100 \text{ MHz}$$

Magnetic field gradient:

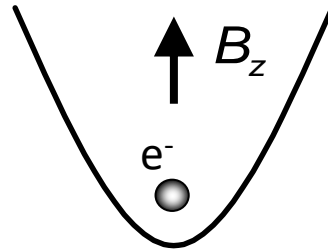
$$\frac{\partial B_x}{\partial z} \approx 0.14 \text{ mT/nm}$$



Decoherence of spin states

SO coupling will decrease coherence of spin states!

① Spin relaxation



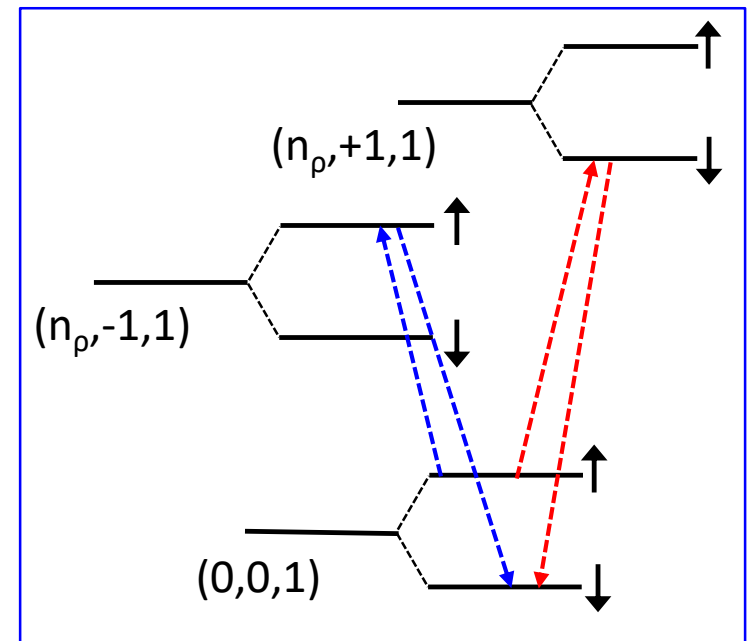
$$H = H_{Fock-Darwin} + H_z + H_s + H_{SO} + H_{e-rippons}$$

$\underbrace{\hspace{10em}}_{\text{in-plane motion}} \quad \underbrace{\hspace{5em}}_{\text{vertical motion}} \quad \underbrace{\hspace{5em}}_{\text{spin}} \quad \underbrace{\hspace{15em}}_{\text{perturbation}}$

Second order perturbation theory:

$$T_1^{-1} \approx \Gamma_{\text{orbital}} \left(\gamma_c I_B \left(\frac{\partial B_\rho}{\partial \rho} \right) \frac{\sqrt{\omega_c^2 + \omega_0^2}}{8\omega_0^2} \right)^2 \approx 10 \text{ s}^{-1}$$

Virtual transitions to orbital states



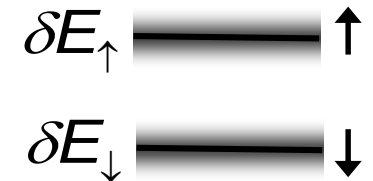
Decoherence of spin states

SO coupling will decrease coherence of spin states!

② Spin dephasing

$$H = \underbrace{H_{\text{Fock-Darwin}}}_{\text{in-plane motion}} + \underbrace{H_z}_{\text{vertical motion}} + \underbrace{H_s}_{\text{spin}} + \underbrace{H_{\text{SO}} + H_{\text{e-ripples}}}_{\text{perturbation}}$$

$$\langle [\varphi(t) - \varphi(0)]^2 \rangle = \frac{1}{\hbar^2} \int_0^t \int_0^{t_2} dt_1 dt_2 \langle \delta E_{\uparrow\downarrow}(t_1) \delta E_{\uparrow\downarrow}(t_2) \rangle$$



Second order perturbation theory:

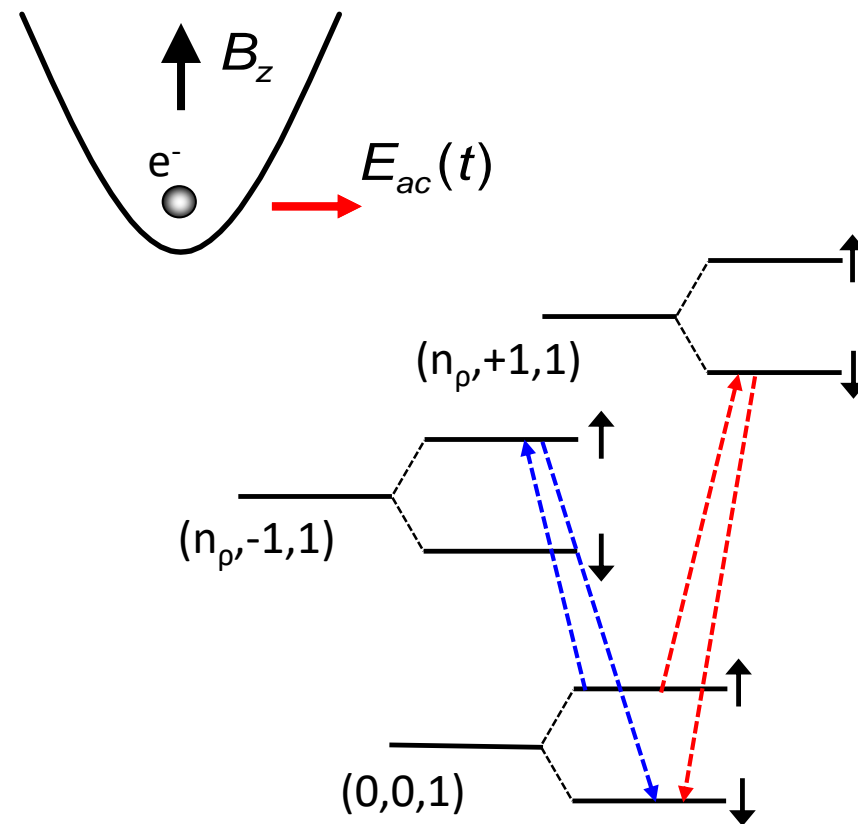
$$\Gamma_{\varphi}(T) \approx 10^{-2} \text{ s}^{-1} \quad @ T=100 \text{ mK}$$



Spin rotation by RF electric field

$$H = H_{Fock-Darwin} + H_z + H_s + H_{SO} + H(t)$$

$\underbrace{\hspace{10em}}_{\text{orbital motion}} \quad \underbrace{\hspace{2em}}_{\text{spin}} \quad \underbrace{\hspace{2em}}_{\text{SO interaction}}$



Electrical dipole approximation:

$$H(t) = V e^{-i\omega t} + V_+ e^{i\omega t}, \quad V = e E_{ac} x$$

Second order perturbation theory:

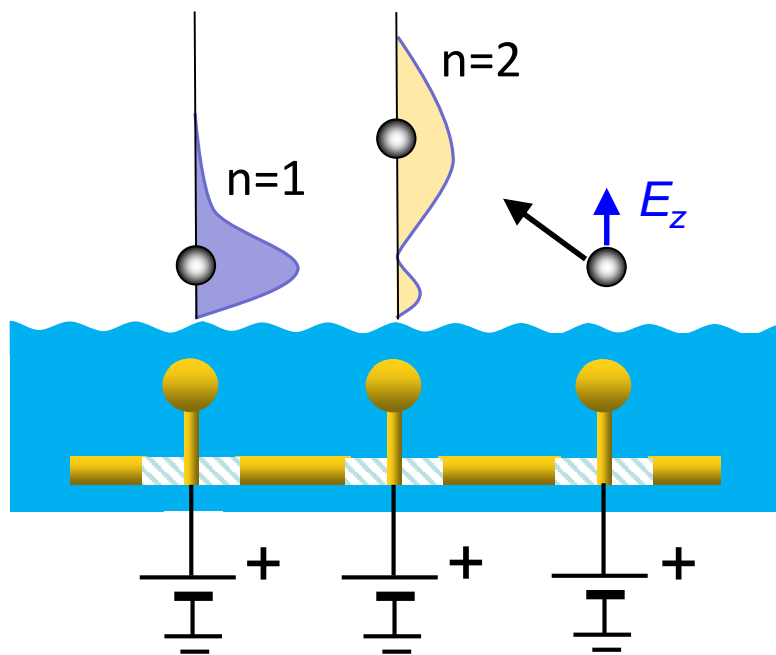
$$\Omega_{rot} \approx \Omega_{dipole} \gamma_c I_B \left(\frac{\partial \mathbf{B}_z}{\partial \rho} \right) \left(\frac{\sqrt{\omega_c^2 + \omega_0^2}}{4\omega_0^2} \right) \approx 20 \text{ MHz} \quad @ E_{ac} = 1 \text{ V/mm}$$



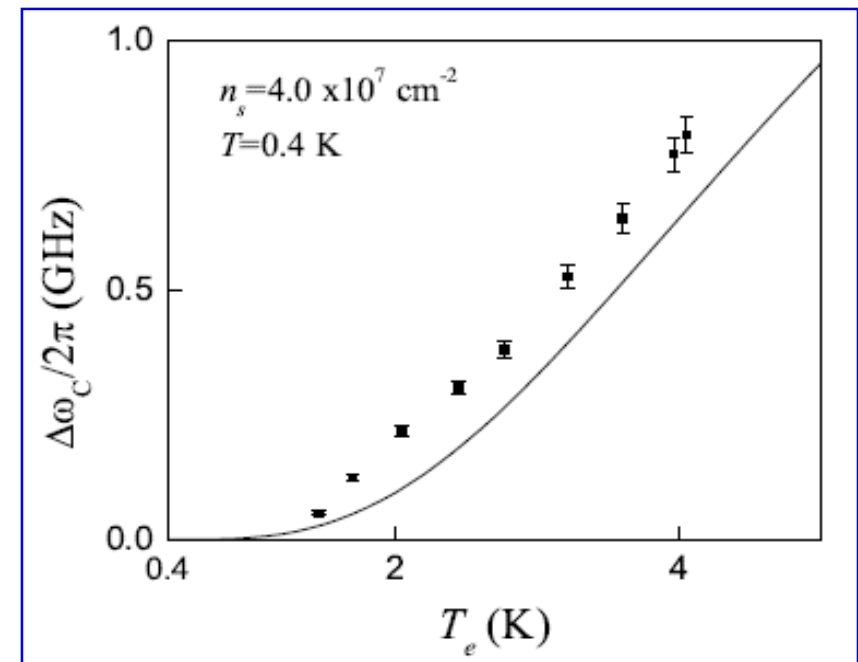
Electron-electron interaction

- Spin relaxation ($T_1=100$ ms)
- Spin dephasing ($T_2^*>1$ s)
- Spin rotations ($\Omega_{\text{rot}}\sim 10$ MHz)

Coulomb shift of Rydberg energies:



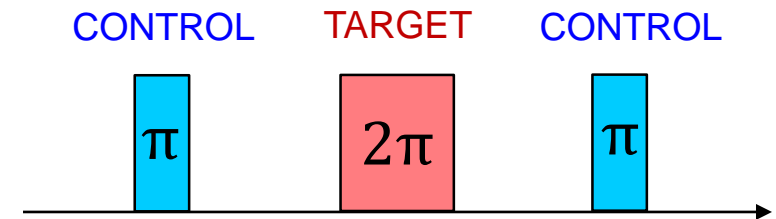
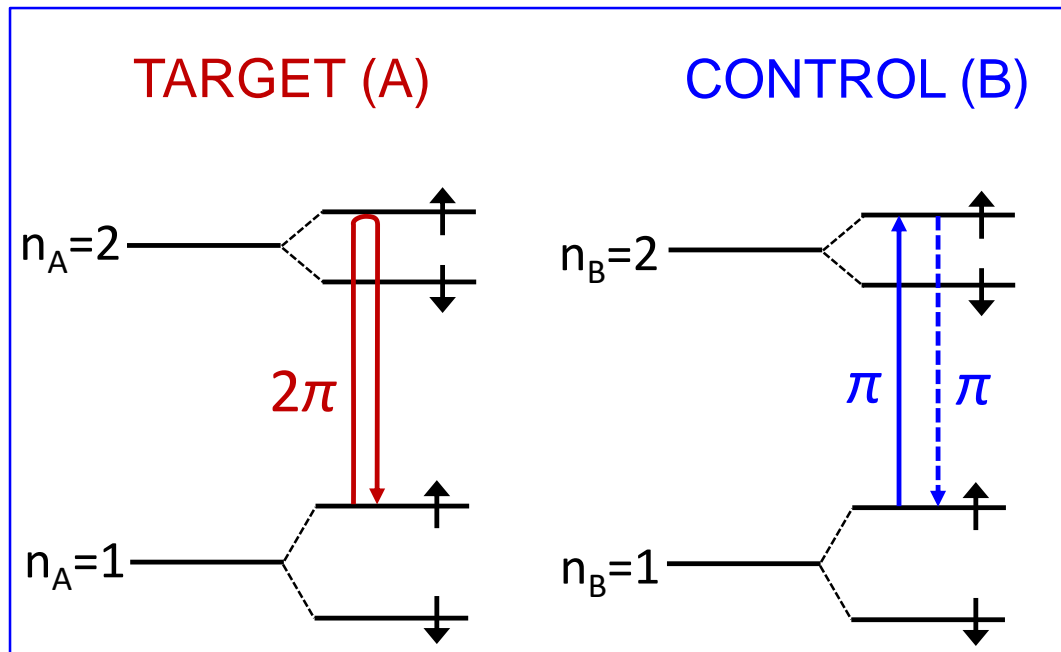
Two-qubit gate?
The Coulomb interaction!



[DK *et al.*, PRL **103**, 096801]



Two-qubit gate



$$\begin{aligned}
 |11\rangle|\downarrow\downarrow\rangle &\rightarrow |11\rangle|\downarrow\downarrow\rangle \rightarrow |11\rangle|\downarrow\downarrow\rangle \rightarrow |11\rangle|\downarrow\downarrow\rangle \\
 |11\rangle|\uparrow\downarrow\rangle &\rightarrow |11\rangle|\uparrow\downarrow\rangle \rightarrow |11\rangle|\uparrow\downarrow\rangle \rightarrow |11\rangle|\uparrow\downarrow\rangle \\
 |11\rangle|\downarrow\uparrow\rangle &\rightarrow -i|12\rangle|\downarrow\uparrow\rangle \rightarrow -i|12\rangle|\downarrow\uparrow\rangle \rightarrow |11\rangle|\downarrow\uparrow\rangle \\
 |11\rangle|\uparrow\uparrow\rangle &\rightarrow -i|12\rangle|\uparrow\uparrow\rangle \rightarrow i|12\rangle|\uparrow\uparrow\rangle \rightarrow -|11\rangle|\uparrow\uparrow\rangle
 \end{aligned}$$

Controlled-phase gate

[Cirac and Zoller, PRL 74, 4091]

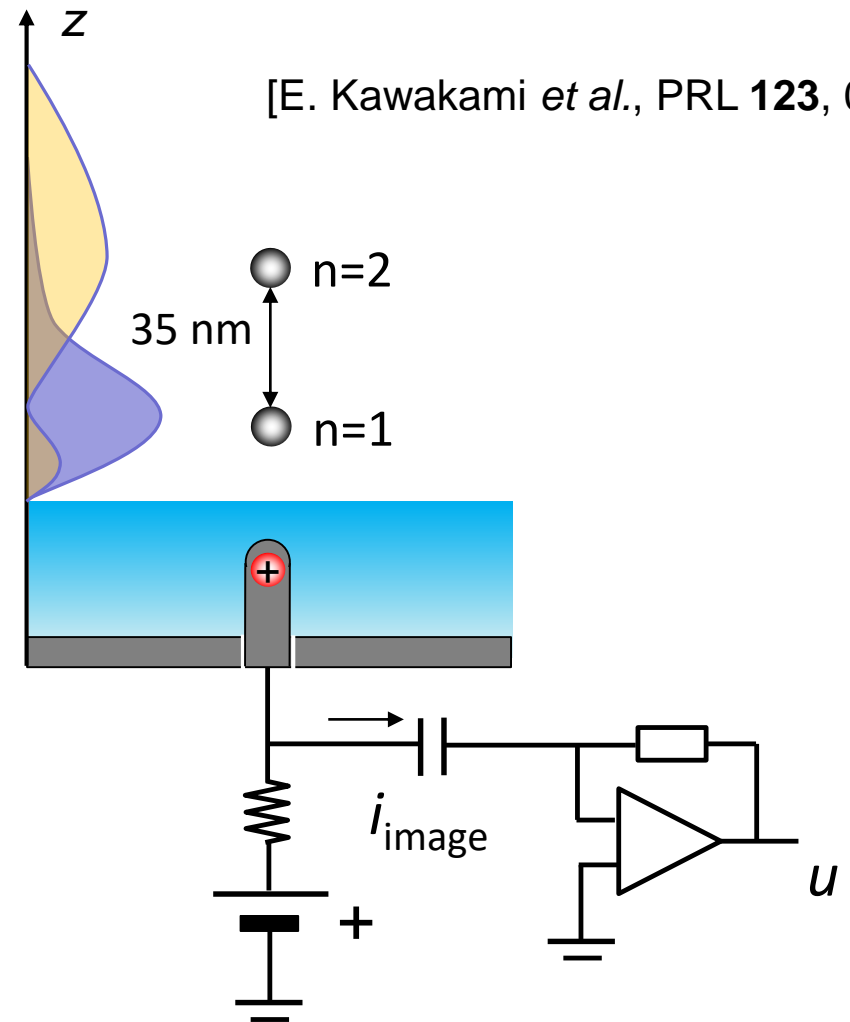
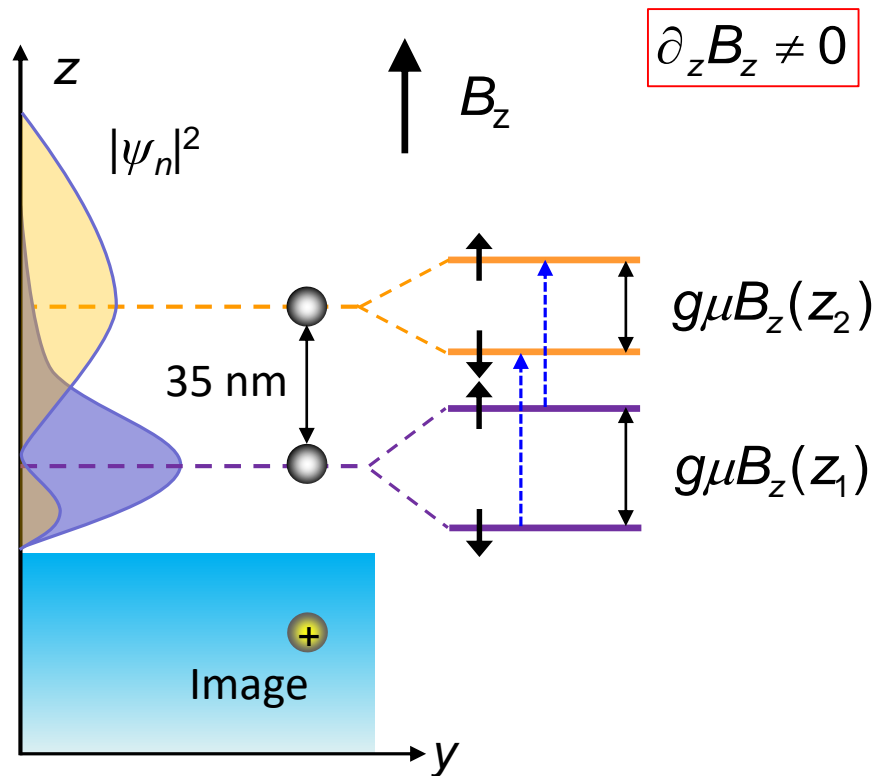
All rotations are on Rydberg states

$$\Omega_{dipole} \sim 100 \text{ MHz}$$

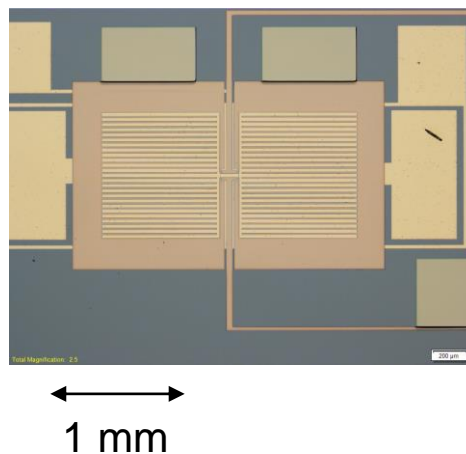
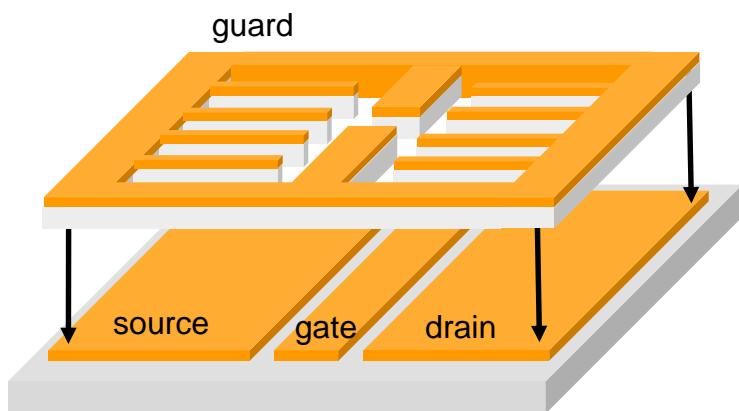


Spin-state readout

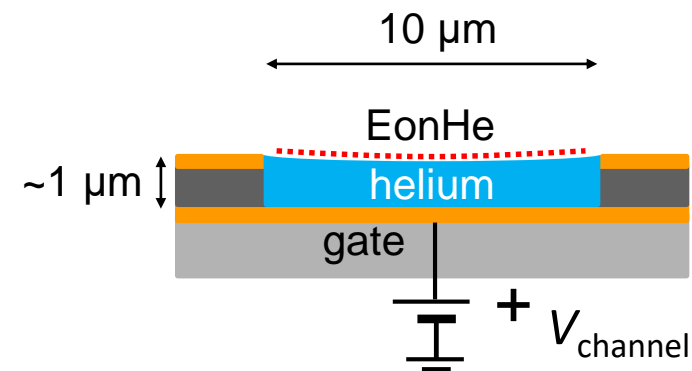
Detect the Rydberg transition $\overset{\text{SO}}{\implies}$ **know the spin state!**



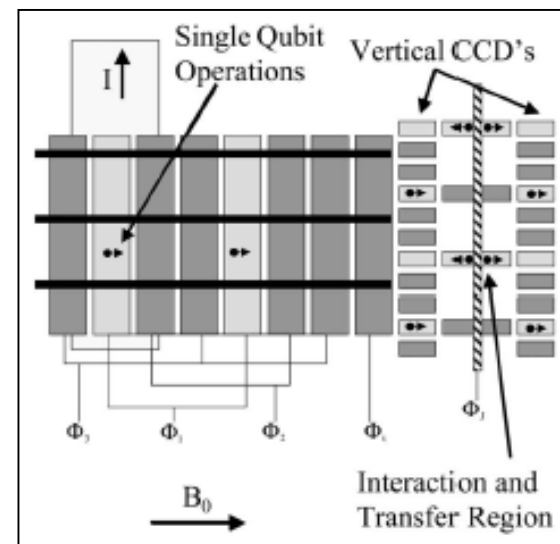
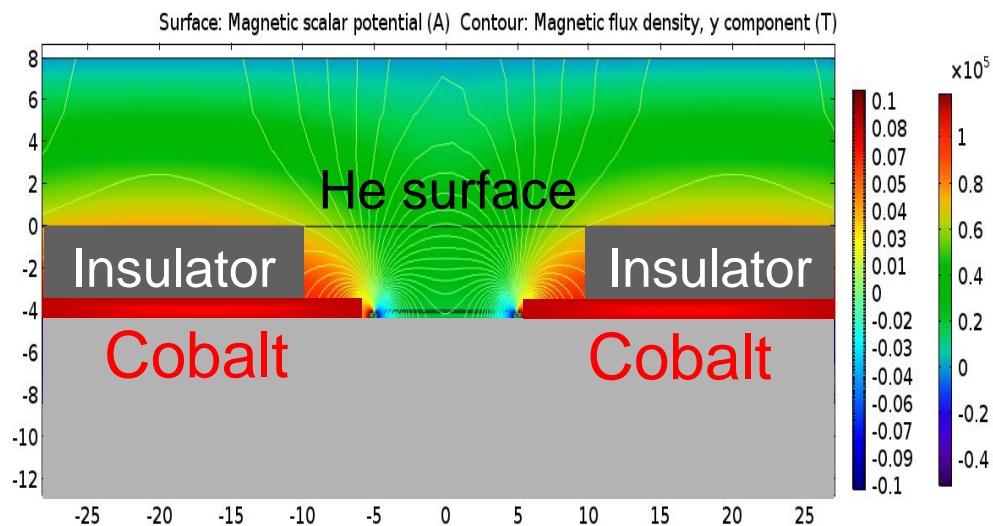
Electrons in Microchannel Devices



Central Channel



B-field gradient:

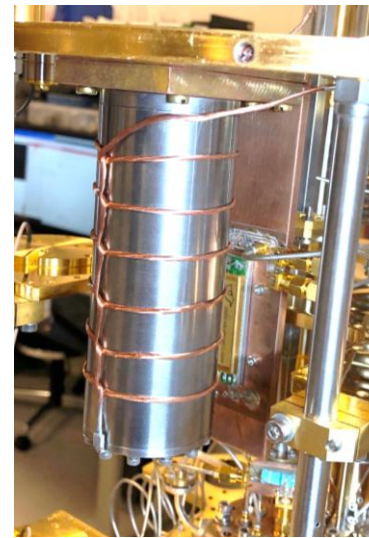
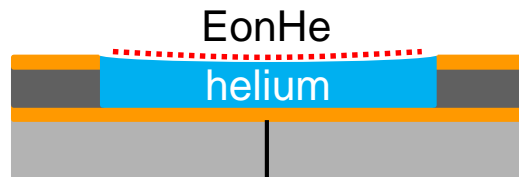
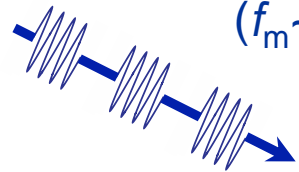


[S. A. Lyon, Phys. Rev. A, **74**, 052338]

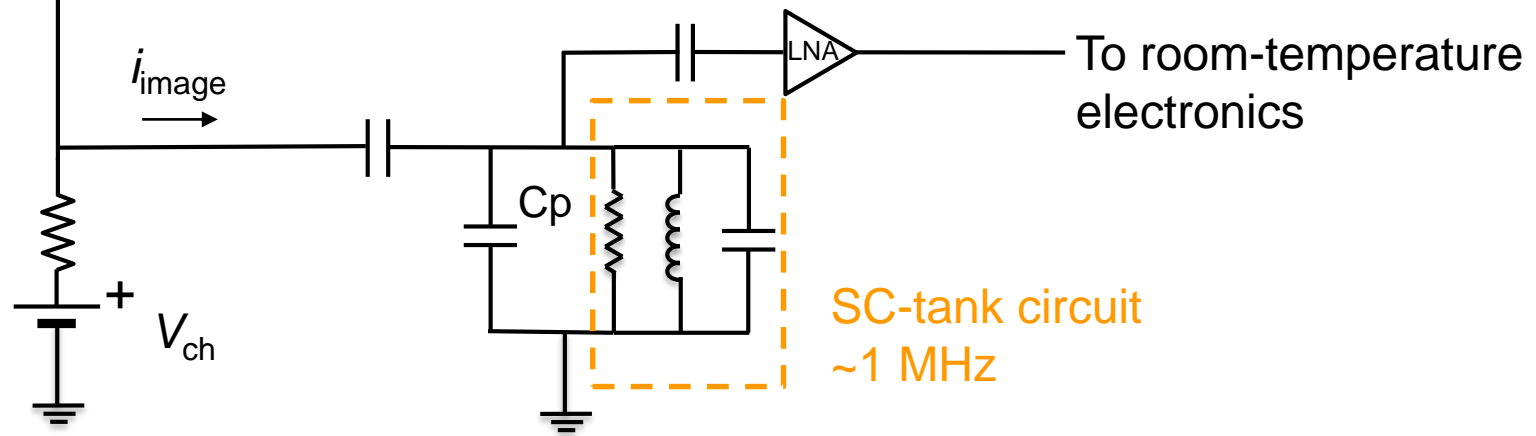
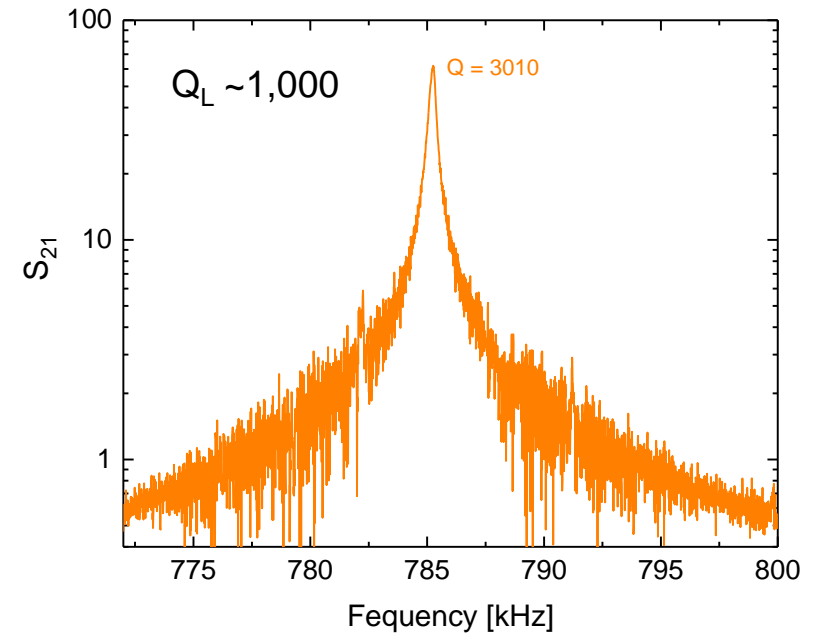


Detection of Rydberg transition of EonHe in a microchannel

Pulse-modulated
100 GHz excitation
($f_m \sim 1$ MHz)



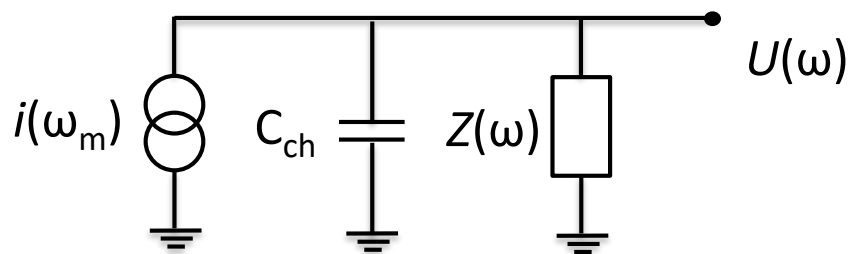
NbTi helical resonator



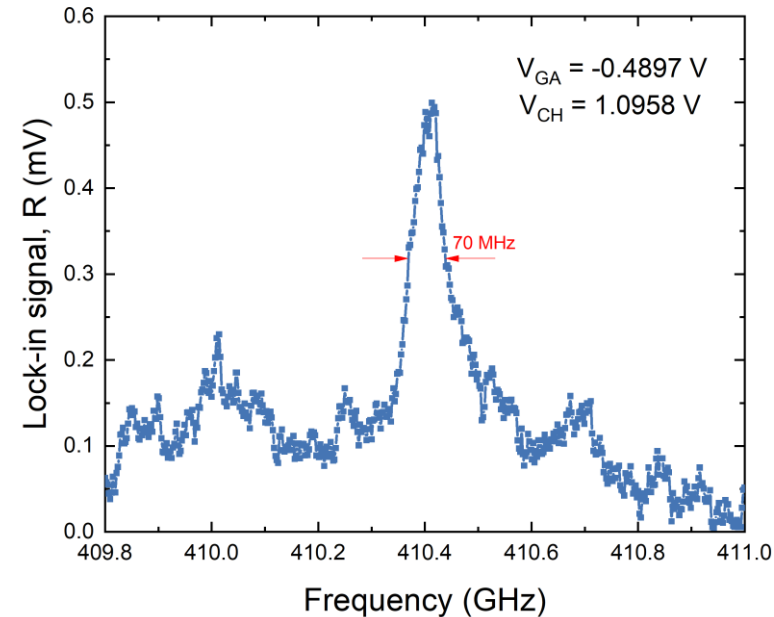
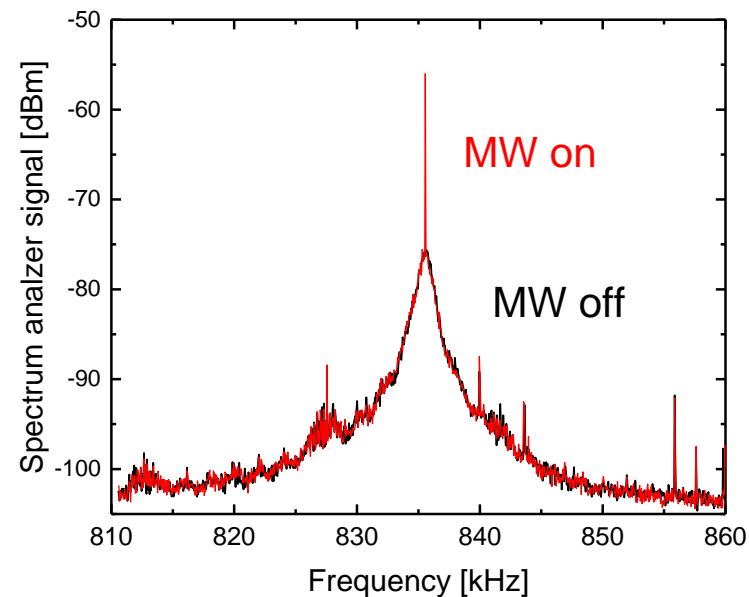
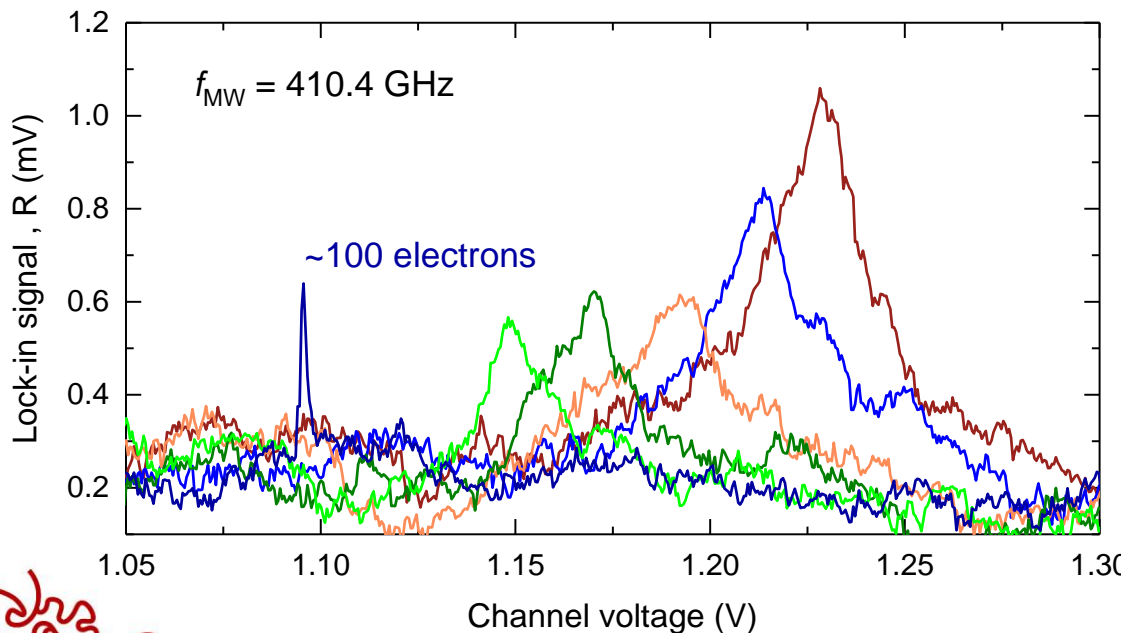
[Wineland and Dehmelt, J. Appl. Phys. **46**, 919]



Rydberg transition



Stark-tuning to resonance:



Summary

Introduce SO coupling to control spin states of EonHe

Suitable for quantum computing

- Slow spin decoherence ($T_1=100$ ms)
- Fast spin rotations ($\Omega_{\text{rot}}\sim 10$ MHz)
- Fast 2-qubit gate ($\Omega_{\text{dipole}}\sim 100$ MHz)
- Spin readout by image current

Microchannel devices: scalability, mobile qubits, QCCD architecture..



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OIST



Mikhail
Belianchikov
(Postdoc)

OIST

<https://www.groups.oist.jp/qdu>



<https://www.oist.jp>

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OIST



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